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# Numerical investigation on the response of ground-reinforced embankments under repeated impact

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ABSTRACT

Ground-reinforced embankment (GRE) is an effective and environmentally friendly technique of rockfall intervention. These earth structures are built with layers of compacted soil alternated with geotextiles, geogrills, metallic wire stripes or nets. GREs are designed to sustain repeated rock impact during their service life, but there is very little experimental or numerical research on the GRE response under such impact conditions. A comprehensive numerical investigation of GRE response under repeated rock impact is carried out. The GRE is built with several layers of sand wrapped by geosynthetics. An advanced elastoplastic constitutive model for sand is adopted. For the GREs built with dense and loose sand, most of the impact energy is dissipated by plastic deformation in the soil. Sand density has a dominant influence on the deformation and failure mechanism of GREs. During repeated impacts, elements near the impact location fail with increasing mean effective stress and Mises stress in dense sand. However, soil elements reach failure as the mean effective stress decreases and the soil has higher stiffness and shear strength. After multiple impacts, shear bands form in loose sand but strain localisation mainly occurs at the impact point for dense sand.

#### 1. Introduction

Rockfall is characterised by rock material's sudden release and movement down a slope or cliff face. This hazard stems from various geological, climatic, and human-induced factors, making it a complex phenomenon to understand and mitigate. Geological factors such as the presence of steep slopes, weak rock formations, and fractures or faults contribute to the instability of rock masses, increasing the likelihood of rockfalls (Peila et al., 2007; Guzzetti et al., 2005; Moos et al., 2022; Bertrand et al., 2013). Weathering processes, including freeze-thaw cycles and rainfall, weaken the cohesion between rocks, facilitating their detachment and subsequent fall (Pérez-Rey et al., 2019). Human activities such as mining, quarrying, and construction can also destabilise rock formations, altering natural slopes and triggering rockfalls (Hungr et al., 2014). Stoffel et al. (2024) have found that continuous global warming is driving more rockfall incidents in the Swiss Alps. Rockfalls pose a direct risk to human life, particularly in areas with high population density or frequented by tourists. Moreover, rockfalls can damage infrastructure such as roads, railways, buildings, and utility

networks, disrupting transportation systems and essential services (Agliardi et al., 2009). The economic consequences of such damage can be substantial, requiring costly repair and maintenance efforts.

Once the mechanism of rockfall (detachment point, rockfall trajectory and deposition of rock debris) and potential impact on infrastructure/human habitats are defined, measures for rockfall protection need to be designed, which include either active or passive ones. The active measures are designed to prevent the detachment of rock from rock slopes, such as rock bolts, rock dowels and rock netting (Bertolo et al., 2009; Peila and Ronco, 2009; Tran et al., 2013; Wyllie, 2015). The passive ones are constructed to intercept and stop the blocks, thus eliminating or reducing the impact of rockfall hazards. Examples of passive measures include rock shelters, catch fences and GREs. Rockfall shelters consist of reinforced concrete structures with roofs covered by a layer of absorbing material, typically soil backfill (Lambert and Bourrier, 2013; Moon et al., 2014; Bourrier et al., 2015; Turner and Schuster, 2012). Rockfall catch fences are constructed with steel wires, posts, ropes and anchors. This steel mesh acts as a barrier to intercept moving blocks and transfers stopping forces to ground anchors. In designing

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#### Table 1

Model parameters for Fontainebleau sand.

Group	Parameter	Value
	ν	0.25
Elasticity	$G_0$	100
	n	0.55
	$e_{\Gamma}$	0.825
Critical state	$\lambda_c$	0.055
Critical state	ξ	0.46
	$\phi_u/^\circ$	33.2
	$k_p$	0.0015
Choor eliding	$A_d$	0.55
Silear-silding	$m_1$	2.5
	$m_2$	2.8
Compression	λ'	0.25
Compression	$p_{c0}/kPa$	15,000
	e <sub>refu</sub>	0.296
Cuoin breekees norematore	e <sub>ref0</sub>	0.825
Gram Dreakage parameters	b/kPa	12,000
	ρ	4

high-energy fences, some energy-dissipation devices can be installed to reduce the forces being transferred to ground anchors (Tran et al., 2013). GREs are used when the rock blocks have volumes or speeds that are great enough to break through the maximum resistance of traditional catch fences (Lambert and Bourrier, 2013; Peila et al., 2007). According to Volkwein and Gerber (2011), the choice of rockfall protection measures depends on their energy absorption capacities. For example, rigid fences are suitable for low-energy impacts up to approximately 100 kJ, while barrier systems can handle impacts ranging from 10 kJ to 10,000 kJ. Retaining walls or GREs, designed for higher energy levels, can withstand impacts ranging from 10 kJ to 100,000 kJ. GREs have two main advantages over the other passive measures: (a) They have a much higher energy absorption capacity and can sustain multiple rockfall impacts (Peila et al., 2007; Lambert and Kister, 2018); (b) They are easier and cheaper to build and repair than other passive measures because local soils can be used in construction and the reinforced materials used are light-weight (Brunet et al., 2009; Xiao et al., 2023; Celik, 2023).

There have been extensive experimental and numerical studies on the response of GREs under rockfall impact, including full-scale tests (e. g., Aminata et al., 2008; Peila et al., 2007; Peila, 2012), small-scale tests (e.g., Kister and Fontana, 2011; Lambert and Kister, 2017; Korini et al., 2021) and numerical modelling (Di Prisco and Vecchiotti, 2006; Shen et al., 2019; Yan et al., 2020; Oggeri et al., 2021; Lu et al., 2021). It is found that numerical modelling is a powerful tool for analysing the response of GRE and assisting the design. The modelling of GREs in this study has been designed based on existing research in which static element test results are used for dynamic modelling (Lu et al., 2021; Gao et al., 2023; Xiao et al., 2023). However, the existing numerical studies have two major limitations. First, the mechanical behaviour of soil in GREs has not been properly accounted for. The Mohr-Coulomb model has been frequently used (Lu et al., 2021; Ma et al., 2024), which is not capable of capturing the highly nonlinear stress-strain relationship of soils (Gao et al., 2014; Liang et al., 2020; Wang et al., 2024; Yao et al., 2004). Secondly, most studies have focused on the response of GREs under a single rockfall impact, while GREs are designed for multiple impacts (Peila et al., 2007; Xiao et al., 2023; Maheshwari et al., 2024; Maheshwari and Bhowmik, 2024). This study aims to use an advanced elastoplastic sand model to simulate the deformation and failure mechanism of GREs under repeated impacts and to provide important theoretical support for the design of GREs in the field.

Numerical modelling of repeated rockfall impacts on GREs is presented in this study. A GRE built with several layers of sand wrapped in geosynthetics is modelled. A double-yield-surface (DYS) sand model accounting for the effect of soil state (density and stress) is used (Jin et al., 2018). The paper is organised as follows. The DYS model and its predictive capability are introduced first to facilitate the result interpretation. Following this, the numerical modelling of repeated rock impact on GREs is presented. The GRE response is analysed in terms of block velocity and energy evolution, impact stress distribution, impact force evolution and deformation localisation. Implications of the findings for GRE design are discussed.

## 2. The Constitutive Model

The yield function for shear sliding  $f_1$  is expressed as (Jin et al., 2018),

$$f_1 = \frac{q}{p} - \frac{M_p \varepsilon_d^p}{k_p + \varepsilon_d^p} = 0 \tag{1}$$

where *q* is the deviatoric stress ( $q = \sqrt{\frac{3}{2}} s_{ij} s_{ji}$ , with  $s_{ij} = \sigma_{ij} - p \delta_{ij}$ ),  $\sigma_{ij}$  is the stress tensor,  $p = \sigma_{ii}/3$  is the mean effective stress,  $\delta_{ij}$  is the Kronecker delta (= 1 for i = j, and = 0 for  $i \neq j$ ),  $\varepsilon_d^p$  is the deviatoric plastic strain,  $k_p$  is a model parameter related to the plastic shear modulus;  $M_p = 6sin \frac{\phi_p}{(3 - sin\phi_p)}$  is the virtual peak stress ratio that varies with the peak friction angle  $\phi_p$  expressed in terms of the critical friction angle  $\phi_u$  and the current void ratio *e*,

$$tan\phi_p = tan\phi_u \left(\frac{e_c}{e}\right)^{m_1} \tag{2}$$

where  $m_1$  is material constant and the critical state void ratio  $e_c$  is expressed as,

$$e_{c} = e_{\Gamma} \exp\left[-\lambda_{c} \left(\frac{p}{p_{a}}\right)^{\xi}\right]$$
(3)



Fig. 1. Comparison between the model simulations and test results of Fontainebleau NE34 sand in drained triaxial compression tests: (a) deviatoric stress versus axial strain, and (b) void ratio versus mean effective stress.



Fig. 2. Model prediction of Fontainebleau NE34 sand in drained triaxial compression tests: (a) and (c) deviatoric stress versus axial strain; and (b) and (d) void ratio versus axial strain.



Fig. 3. The cross-section of the model embankment.

where  $p_a$  is atmospheric pressure (101 kPa),  $\lambda_c$  and  $\xi$  are material parameters and  $e_{\Gamma}$  is a variable that evolves with particle breakage which will be given below. The partial differential equations for the plastic potential surface  $g_1$  is given by,

$$\frac{\partial g_1}{\partial \sigma_{ij}} = \frac{\partial g_1}{\partial p} \frac{\partial p}{\partial \sigma_{ij}} + \frac{\partial g_1}{\partial q} \frac{\partial q}{\partial \sigma_{ij}} = A_d \left( M_{pt} - \frac{q}{p} \right) \frac{\partial p}{\partial \sigma_{ij}} + \frac{\partial q}{\partial \sigma_{ij}}$$
(4)

where  $A_d$  is a model parameter,  $M_{pt} = 6 \sin \phi_{pt} / (3 - \sin \phi_{pt})$  with  $\phi_{pt}$  being expressed as,

$$tan\phi_{pt} = tan\phi_u \left(\frac{e_c}{e}\right)^{-m_2} \tag{5}$$

where  $m_2$  is a model parameter. Note that the explicit formulation for  $g_1$  cannot be derived based on Eqs. (4) and (5). The second yield function  $f_2$  and the plastic potential surface  $g_2$  of the model which is used to describe the normal compression behaviour of sand are expressed as,



Fig. 4. Mesh size of GRE used in the simulations.

Table	2
Tubic	-

Properties	for	the	block,	geosynthetic and	sand.

	Density (kg/ $m^3$ )	Young's modulus (MPa)	Poisson's ratio $\nu$
Block Geosynthetic Sand	7830 1400 1630	/ 180 /	/ 0.25

$$f_2 = g_2 = \frac{1}{2} \left( \frac{q}{M_p p} \right)^3 p + p - p_c$$
 (6)

Table 3Summary of impact tests.

Simulation		Impact velocity	Impact energy	Accumulated energy
numbers		(m/s)	(kJ)	(kJ)
$e_0 = 0.6$	$e_0 = 0.8$			
1-1	2–1	10	44.25	44.25
1–2	2–2	10	44.25	88.50
1–3	2–3	10	44.25	132.75

where  $p_c$  is the hardening parameter controlling the size of the yield surface. The yield surface expands with the increment of plastic volumetric strain increment  $de_{p}^{p}$ ,

$$dp_c = p_c \frac{1+e}{(\lambda'-G')e} d\varepsilon_{\nu}^p \tag{7}$$

where the parameter  $\lambda'$  controls the slope of the compression line, G' = p(1 + e)/G is a normalised value calculated by elastic bulk modulus *G* which is defined as,

$$G = G_0 p_a \frac{(2.97 - e)^2}{(1 + e)} \left(\frac{p}{p_a}\right)^n$$
(8)

where  $G_0$  and n are elastic model parameters. The particle breakage is modelled by considering the evolution of  $e_{\Gamma}$  in Eq. (3) with the grain breakage index  $B_r^*$ 

$$\boldsymbol{e}_{\Gamma} = \boldsymbol{e}_{refu} + \left(\boldsymbol{e}_{ref0} - \boldsymbol{e}_{refu}\right) \boldsymbol{e} \boldsymbol{x} \boldsymbol{p} \left(-\rho \boldsymbol{B}_{r}^{*}\right) \tag{9}$$

where  $e_{ref0}$  and  $e_{refu}$  are the virgin and ultimate reference critical void ratios for the virgin soil without grain breakage,  $\rho$  is the material parameter and  $B_r^*$  is expressed as,



Fig. 5. Evolution of block velocity during the first and third impacts: (a) initial void ratio of sand  $e_0 = 0.6$ ; (b) initial void ratio of sand  $e_0 = 0.8$ .



Fig. 6. Penetration depth of the block in GREs different void ratios.



**Fig. 7.** Energy evolution in the first impact: (a) initial void ratio of sand  $e_0 = 0.6$ ; (b) initial void ratio of sand  $e_0 = 0.8$ .



Fig. 8. Energy evolution in the third impact: (a) initial void ratio of sand  $e_0 = 0.6$ ; (b) initial void ratio of sand  $e_0 = 0.8$ .



Fig. 9. Effect of repeated impacts on the energy evolution in GREs with different void ratios: (a) strain energy; (b) plastic dissipation; (c) Frictional dissipation.

$$B_r^* = \frac{w_p}{b + w_p} \tag{10}$$

where  $w_p = \int (p \langle d\varepsilon_v^p \rangle + q d\varepsilon_d^p)$  is the plastic work, are the McCauley brackets and *b* is a material constant controlling the evolution rate of the grain breakage index.

The increment of total strain  $d\varepsilon_{ij}$  is composed of the elastic and plastic strain parts,

$$d\varepsilon_{ii} = d\varepsilon_{ii}^e + d\varepsilon_{ii}^p \tag{11}$$

The elastic and plastic strain increment tensor can be expressed as,

$$d\varepsilon_{ij}^{e} = \frac{1+\nu}{E} d\sigma_{ij} - \frac{\nu}{E} d\sigma_{kk} \delta_{ij}$$
(12)

$$d\varepsilon_{ij}^{p} = \left(d\varepsilon_{ij}^{p}\right)_{1} + \left(d\varepsilon_{ij}^{p}\right)_{2} = d\lambda_{1}\frac{\partial g_{1}}{d\sigma_{ij}} + d\lambda_{2}\frac{\partial g_{2}}{d\sigma_{ij}}$$
(13)

where  $\nu$  and  $E = 3G(1 - 2\nu)$  are Poisson's ratio and Young's modulus,  $d\lambda_1$  and  $d\lambda_2$  are the plastic multipliers representing the shear sliding and compression components, respectively. Derivation of the constitutive equations can be found in Jin et al. (2018).

This model has been implemented in the finite element package ABAQUS/Explicit through the user-material (VUMAT) (Dassault Systèmes, 2023). The cutting plane algorithm of the DYS model is employed in the model implementation (Jin et al., 2018) which is a semi-implicit scheme that can guarantee accurate and efficient results even for large time steps. The model parameters for Fontainebleau sand are given in Table 1. The model predictions of Fontainebleau sand



Fig. 10. Strain energy distribution in geosynthetics with  $e_0 = 0.6$  after the first impact (a, b) the third impact (c, d).

behaviour in drained triaxial compression tests are shown in Figs. 1 and 2 (Jin et al., 2018). The test data is from Andria-Ntoanina et al. (2010). More validation results can be found in Jin et al. (2018).

## 3. Finite element modelling of rockfall impact on GRE

The GRE modelled here is built using 8 layers of sand wrapped by geosynthetics. The dimension and mesh size of the GRE used in simulations are shown in Fig. 3 and Fig. 4, respectively. The slope angle is chosen based on the suggested values reported in the literature (Peila et al., 2007; Ronco et al., 2009). The properties of the block, geosynthetics and sand are given in Table 2.

The block density specified for the block in Table 2 is based on the assumption that the block is made of steel. The pressure-dependent elastic parameters for sand are given in Table 1. Sand is modelled using the eight-node brick element with reduced integration (C3D8R). The geosynthetics with a thickness of 0.25 mm are modelled using 4node quadrilateral membrane elements with reduced integration (M3D4R). The thickness and the mechanical parameters of the geosynthetics used here are the same as those in the Lu et al. (2021). The friction coefficient for contacting geosynthetic surfaces is 0.46 (Ronco et al., 2009; Oggeri et al., 2021). The steel block with a diameter of 600 mm is simulated using rigid 3D bilinear quadrilateral elements (R3D4). Steel is used following the tests by Lu et al. (2021). The friction coefficient between the block and geosynthetics is 0.45 (Uesugi et al., 1989). The bottom of GRE is fully fixed while all the other sides are free to deform. Following the definitions in existing studies (Ronco et al., 2009; Lambert and Bourrier, 2013; Oggeri et al., 2021), the impact side of the GRE is called the mountain side and the opposite one is called the valley side (Fig. 4).

Since the embankment has three free surfaces with two inclined, the

conventional  $K_0$  procedure cannot be used to generate the initial stress state. Therefore, the gravitational loading method is adopted (Gao et al., 2021). Specifically, the embankment is first assumed to be an elastic material with a large elastic modulus and Poisson's ratio  $\nu_e$ . The  $\nu_e$  can be calculated from a lateral earth pressure coefficient  $K_0 = \nu_e/(1 - \nu_e) =$ 0.48 (Jin et al., 2018). The gravitational body force is then applied to the embankment to generate the initial stress state. During the impact modelling, the elastoplastic model with parameters listed in Table 1 is used. The impact tests simulated are summarised in Table 3. In the subsequent impacts, the block is first removed and the GRE is allowed to deform under gravity until the system reaches equilibrium. The equilibrium stress state is then imported to the system as the initial stress state before impact. The initial void ratio of sand is  $e_0 = 0.6$  ( $D_r = 83\%$ ) and  $e_0 = 0.8$  ( $D_r = 21\%$ ) in the simulations (Table 3). It should be mentioned that loose soil is not recommended in the practical design of GREs, loose sand is chosen here to illustrate the effect of density on the GRE response.

GREs built with coarse sand have higher energy absorption capacity as the sand strength and stiffness are higher (Lu et al., 2021). Since there is no triaxial test data on coarse sand for getting the model parameters, Fontainebleau sand with fine particles is used in this study. Consequently, the impact energy used in this study is relatively low (Table 3). However, the findings are expected to be valid for GREs built with coarse sand.

## 3.1. Block velocity and energy evolution under repeated impacts

Fig. 5 shows the evolution of block velocity in the first and third impact for GREs with different initial void ratios. The block bounces back after the first impact, and the velocity increase after impact (at about 0.1 s) is driven by the free fall of the block. In the following two



Fig. 11. Strain energy distribution in geosynthetics with  $e_0 = 0.8$  after the first impact (a, b) the third impact (c, d).



Fig. 12. Plastic dissipation distribution in GREs with  $e_0 = 0.6$  after the first impact (a and b) and the third impact (c and d).

impacts, the block gets embedded in the GRE and the velocity reaches 0 at the end. The penetration depth of the block is higher in loose sand. It increases with the number of impacts (Fig. 6).

During the impacts, part of or all the impact energy (initial kinetic energy of the block) is converted into plastic dissipation in the sand, friction dissipation (friction between geosynthetics and between geosynthetics and block) and strain energy in both sand and geosynthetics (Figs. 7 and 8). It is found that over 74 % of the impact energy is dissipated due to plastic deformation in sand. Frictional dissipation and strain energy account for about 6–10 % and 10–15 % of the impact energy, respectively. This is in agreement with the previous studies (Ronco et al., 2009; Lambert and Bourrier, 2013; Oggeri et al., 2021). In the second and third impacts, the block becomes embedded within the

GRE, resulting in higher percentages of plastic dissipation, frictional dissipation, and strain energy for both loose and dense sand (Fig. 9). More plastic dissipation and strain energy occurs in loose sand due to more deformation in the soil. Meanwhile, frictional dissipation in loose sand is very small (Fig. 9). This indicates that GREs built with less dense sand are more effective in dissipating the impact energy. However, as will be shown later, there is much more deformation in GREs with loose sand.

Fig. 10 and Fig. 11 show the effect of sand void ratio on strain energy in geosynthetics. After the first impact, strain energy in the geosynthetics with dense sand appears in the fourth and fifth layers of geosynthetics on the mountain-side, with a small amount of strain energy also appearing in the seventh and eighth layers near the bottom of



Fig. 13. Plastic dissipation distribution in GREs with  $e_0 = 0.8$  after the first impact (a and b) and the third impact (c and d).



Fig. 14. Effect of repeated impacts on the Mises stress in GREs: (a and b)  $e_0 = 0.6$ ; (c and d)  $e_0 = 0.8$ .

the GRE (Fig. 10a). For the GRE built with loose sand, strain energy appears only in the fourth, fifth, and sixth layers of geosynthetic (Fig. 11a). In the side view of geosynthetics changes shown in Fig. 11 (d), significant deformation appears in the fourth and fifth layers of geosynthetics for the loose sand when compared to the dense sand in Fig. 10 (d), and strain energy also appears on the valley-side of the GRE. This is due to that the GRE built with loose sand undergoes greater deformation after being impacted by a block. Additionally, after the first impact, the magnitude of strain energy is greater in the GRE with dense sand (Fig. 10 and Fig. 11a). But after the third impact, the magnitude of strain energy is greater in the GRE with loose sand (Fig. 10c and Fig. 11c). Therefore, under repeated impact conditions, the efficiency of

geosynthetics in the GRE with loose sand is significantly improved, allowing it to absorb more impact energy.

Figs. 12 and 13 show that the magnitude of plastic dissipation is greater for the GRE with dense sand regardless of whether it is the first or the third impact compared to the GRE with loose sand. For GRE with dense sand, the region of plastic dissipation is mainly concentrated in the vicinity of the block impact area (Fig. 12). However, as the sand void ratio increases, the region of plastic dissipation not only concentrates near the block impact area but also extends towards the toe of the slopes on the mountain-side and valley-side of the GRE (Fig. 13). After the third impact (Fig. 13d), significant plastic dissipation appears at the toe of the slope on the mountain-side. Therefore, repeated impacts have a greater



Fig. 15. Effect of repeated impacts on pressure in GREs: (a and b)  $e_0 = 0.6$ ; (c and d)  $e_0 = 0.8$ .



Fig. 16. Effect of repeated impacts on impact force in GREs: (a)  $e_0 = 0.6$ ; (b)  $e_0 = 0.8$ .



Fig. 17. Effect of repeated impacts on the crest deformation of GREs with varied void ratio: (a)  $e_0 = 0.6$  and (b)  $e_0 = 0.8$ .



Fig. 18. Effect of repeated impacts on the Horizontal deformation of GREs with varied void ratio: (a)  $e_0 = 0.6$  and (b)  $e_0 = 0.8$ .



Fig. 19. Equivalent plastic strain distribution in GRE with  $e_0 = 0.6$  after (a) the first impact and (b) the third impact.

effect on the toes of the GRE when sand density is larger.

## 3.2. Impact stress and impact force under repeated impacts

Figs. 14 and 15 show the Mises stress and pressure distribution in the middle of the GRE before and after impact, respectively. Both the Mises stress and pressure increase after each impact, and their maximum values are observed near the impact point. In dense sand, the Mises stress and pressure after impact increase with the impact numbers. This indicates that the soil elements fail due to an increase in both mean

effective stress and Mises stress. However, the soil elements fail with increasing Mises stress and decreasing pressure under repeated impacts in loose sand (Figs. 14 and 15). In GREs constructed with loose sand, repeated impacts gradually create a crater at the impact point. This crater expands with each subsequent impact. The pressure in the sand is inversely proportional to the crater area. In loose sand, the crater area increases significantly after each impact, leading to lower pressure. In dense sand, the crater area shows much less change, and therefore, the pressure increases with the impact number.

Fig. 16 shows the evolution of horizontal impact force (in the x-



Fig. 20. Equivalent plastic strain distribution in GRE with  $e_0 = 0.8$  after (a) the first impact and (b) the third impact.



Fig. 21. Void ratio distribution in GRE with  $e_0 = 0.6$  after (a) the first impact and (b) the third impact.



Fig. 22. Void ratio distribution in GRE with  $e_0 = 0.8$  after (a) the first impact and (b) the third impact.

direction) on the GRE with different sand densities. The impact force is obtained by summing the horizontal reaction force of all nodes at the bottom of the GRE. The maximum impact force shows very little change in dense sand but increases with the number of impacts in loose sand. In the second and third impacts, the impact force before 0.03 s is negative, which is different from that in the first impact. The main reason is that the impact area of the GRE has been damaged during the first impact. In the subsequent impacts, the impact area of the GRE first tends to move

towards the block, resulting in a negative impact force measured at the bottom.

#### 3.3. Deformation and strain localisation in GREs

Fig. 17 shows the crest deformation of GREs after repeated impacts. The positive vertical displacement indicates heaving, while the negative one signifies settlement. When  $e_0 = 0.6$ , the closer to the mountain-side



**Fig. 23.** Grain breakage index  $B_r^*$  distribution in GRE with  $e_0 = 0.6$  after (a) the first impact and (b) the third impact.



**Fig. 24.** Grain breakage index  $B_r^*$  distribution in GRE with  $e_0 = 0.8$  after (a) the first impact and (b) the third impact.

face (impact side), the greater the heaving, which decreases with the number of impacts. Repeated impacts cause heaving accumulation in the middle part of the crest. When  $e_0 = 0.8$ , the crest near the mountainside shows a slight settlement after the first impact and then experiences slight heaving after the third impact. The settlement reaches the maximum at the middle of the crest and the minimum at the valley side. Repeated impacts cause an increase in the settlement (Fig. 17b).

Fig. 18 illustrates the horizontal deformation of GREs under repeated impacts for two initial void ratios ( $e_0 = 0.6$  and  $e_0 = 0.8$ ). The horizontal deformation is measured from the middle part of the GREs, from bottom to top. For  $e_0 = 0.6$  (Fig. 18a), the maximum horizontal deformation is observed near the top of the GRE, with deformation increasing rapidly along the embankment height. This indicates a localised response with higher stiffness and resistance near the base. In contrast, for  $e_0 = 0.8$  (Fig. 18b), the deformation is more widely distributed, extending higher along the embankment. Although loose sand exhibits greater overall deformation due to its lower stiffness and higher compressibility, it effectively distributes stresses across a broader area. Both cases demonstrate increasing deformation with repeated impacts, but the rate of increase is more pronounced in loose sand, particularly near the embankment's base. These findings suggest that GREs constructed with dense sand provide higher localised resistance.

Figs. 19 and 20 show the equivalent plastic strain (EPS) distribution in GREs with different  $e_0$ . To avoid the impact of geosynthetics on the observation of EPS changes in sand, they are not shown in these figures. As expected, the EPS is larger in the GRE with a bigger  $e_0$ . Strain localisation occurs around the impact area after the first impact. After the third impact, the strain localisation extends to the top of GREs built with both dense and loose sand. A clear shear band can be observed in loose sand after the third impact (Fig. 20). In dense sand, a shear band has just started to form at the end of the third impact, and a fully developed one is impacted after more impacts.

The void ratio distribution after impact is shown in Figs. 21 and 22. After the first impact, there is a significant increase in the void ratio (volume expansion) in the vicinity of the impact point. However, this part becomes denser after the third impact. The void ratio distribution in GRE varies with  $e_0$ . In dense sand with  $e_0 = 0.6$ , the top, bottom, and the area closer to the impact point of the GRE show more volume expansion. The area of volume expansion spreads towards the valley side after the third impact. Such volume expansion causes the heaving of GRE at the crest as shown in Fig. 21. In loose sand with  $e_0 = 0.8$ , volume expansion mainly occurs in the middle of the GRE. Volume contraction is observed at the top of the GRE built with loose sand.

Figs. 23 and 24 illustrate the impact of different  $e_0$  on the distribution of the Grain Breakage Index  $(B_r^*)$ . Positive  $B_r^*$  indicates the occurrence of grain breakage, with higher values signifying a greater degree of breakage. In dense sand, the degree of breakage on the GRE surface is higher, but the range of breakage is smaller. Conversely, in loose sand, the degree of breakage is lower, but the range of breakage is wider. However, it should be emphasised that the particle breakage is not significant in these tests because high stress occurs only near the impact area.

#### 4. Conclusions

The response of GREs built with sand wrapped by geosynthetics is investigated using finite element modelling. A DYS model that accounts for the effect of density and pressure on sand response is used. GREs with two different void ratios are simulated. A GRE built with very loose sand is modelled to investigate the effect of sand density on the GRE response. The main conclusions are:

- 1. The impact energy is converted to plastic dissipation in sand, friction dissipation and strain energy in both sand and geosynthetics. There is more plastic dissipation in the GRE with loose sand due to more plastic strain accumulation.
- Under repeated impacts, the soil elements in dense sand reach failure with increasing Mises stress and mean effective stress. However, loose sand reaches failure as the mean effective stress decreases and Mises stress increases.
- 3. Strain localisation appears around the impact area after the first impact. A clear failure surface can be observed after the third impact in loose sand, which starts at the impact point and ends on the top of the GRE. In dense sand, a fully developed shear band is expected after more impacts.
- 4. There is heaving on the top of GRE with dense sand which is caused by volume expansion in the soil after impact. In the GRE with loose sand, slight heaving occurs near the mountain side and settlement dominates in the remaining par, which is due to the volume contraction of sand at the top and bottom parts of the GRE.

This study has shown that GRE built with very loose sand can also sustain multiple impacts without collapsing. There are two main reasons for this. First, loose sand can dissipate a significant amount of impact energy due to the plastic strain accumulation. Secondly, the sand has been wrapped by geosynthetics which strengthens the soil and prevents collapsing. In practical design, dense sand is still preferred as it shows less deformation accumulation after repeated impacts. However strict density control may not be very important when geosynthetics with sufficient stiffness and strength are used, which will be investigated in the future by small-scale and full-scale field tests.

### CRediT authorship contribution statement

Xin Li: Writing – review & editing, Writing – original draft, Visualization, Methodology, Investigation, Conceptualization. Zhiwei Gao: Writing – review & editing, Writing – original draft, Methodology, Investigation, Funding acquisition, Formal analysis, Conceptualization. Liang Lu: Methodology, Funding acquisition, Conceptualization.

### Declaration of competing interest

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

Dr. Zhiwei Gao reports financial support was provided by The Royal Society. If there are other authors, they declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Data availability

The research data presented in the present paper is available upon request.

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