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# Evaluation of three weight functions for nonlocal regularisation of sand models

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**Abstract:** Nonlocal regularisation is frequently used to resolve the mesh-dependency issue caused by strain softening in finite element (FE) simulations. Some or all the variables affecting strain softening are assumed to depend on the local and/or neighbouring ones in this method. The weight function is the most component of a regularisation method. There are three most widely used weight functions, including the Gaussian distribution (GD), Galavi and Schweiger (G&S) and over-nonlocal (ON) functions. Though all of them are found to alleviate or eliminate the mesh dependency in simple boundary value problems (BVPs) like plane strain compression, evaluation of their performance in real-world BVPs is rare. A detailed comparison of these functions has been carried out based on an anisotropic sand model accounting for the evolution of anisotropy. The increment of void ratio is assumed nonlocal. All functions give mesh-independent force-displacement relationship in drained and undrained plane strain compression tests. The shear band thickness shows a small variation when the mesh size is smaller than the internal length. None of them can eliminate the mesh dependency of shear band orientation. The G&S method is the most efficient in eliminating the mesh dependency in the strip footing problem. The ON method can give excessive overprediction of volume expansion around strip footings, leading to unrealistic low reaction forces on strip footings at large deformation. All three weight functions give mesh-independent results for the earth pressure acting on a retaining wall.

**Keywords:** Sand, constitutive model, strain softening, nonlocal regularisation, weight function

## 33 **1. Introduction**

34 The response of the FE continuum after it has reached its peak is inherently dependent on  
35 the mesh used, as noted by Bazant and Jirasek (2002). For strain localisation analysis, the FE  
36 solution will converge to a unique one as the mesh size gets smaller, when a strain-hardening  
37 model is used (e.g., Galavi and Schweiger, 2010; Lu, et al., 2019). But such convergence cannot  
38 be obtained when a strain-softening model is used (e.g., Mallikarachchi and Soga, 2020; Gao  
39 et al., 2022). Mathematically, this mesh dependency is linked to the transformation of the  
40 governing partial differential equations from elliptic to hyperbolic, which occurs when the  
41 material behaviour transitions from hardening to softening. Previous research has also  
42 highlighted this issue (e.g., Mühlhaus, 1986; Galavi and Schweiger, 2010; Guo and Stolle, 2013;  
43 Lu, et al., 2019; Cui et al., 2023). Alsaleh et al. (2006) have pointed out that the FE simulation  
44 of strain localization is subject to mesh dependency due to the use of classical continuum  
45 models that do not account for micro-structural factors, such as particle size and associated  
46 voids.

47

48 Different methods have been proposed to resolve the mesh-dependency issue, including the  
49 strain-gradient enhanced approaches (e.g., Aifantis, 1984; de Borst and Mühlhaus, 1992;  
50 Dorgan and Voyiadjis, 2003), micro-polar plasticity approach (e.g., Mühlhaus, 1986; Alshibli  
51 et al., 2006; Tang et al., 2013), nonlocal regularisation method (e.g., Eringen, 1972; Lü et al.,  
52 2009; Galavi and Schweiger, 2010; Guo and Stolle, 2013; Lazari et al., 2015; Summersgill et  
53 al., 2017; Mallikarachchi and Soga, 2020; Singh et al., 2021; Gao et al., 2022; Cui et al., 2023)  
54 and viscous plasticity theory (e.g., Oka, et al., 1995; Wang, et al. 1997; Higo, 2004; Yin et al.  
55 2010). An internal length scale is introduced to the constitutive model formulation in these  
56 methods, which controls the degree of deformation localisation and preserves the well-  
57 posedness of the governing partial differential equations irrespective of the refinement of  
58 the mesh (de Borst et al., 1993). Among these methods, nonlocal regularisation is the most  
59 widely used for advanced soil models. Nonlocal methods are proposed based on the  
60 hypothesis that the response of materials depends on the deformation field of a local  
61 material point and a weighted average of its neighbouring points (Mallikarachchi, 2019).

62

63 The weight function is the most important component of a nonlocal regularisation method.  
64 The GD function has been used in many early studies (Eringen, 1974; Bažant et al., 1984). The

65 variable at the current stress point contributes most to the nonlocal one, and therefore, the  
66 nonlocal variable is concentrated at the local point and cannot spread to surrounding points.  
67 Galavi and Schweiger (2010) proposed a new weight function in which the local variable does  
68 not affect the nonlocal one. Moreover, Vermeer and Brinkgreve (1994) have proposed the  
69 over-nonlocal method which uses a linear combination of the local and the nonlocal variables.  
70 A nonlocal parameter  $m$  is introduced to control the proportion of local and nonlocal  
71 variables in weight functions. Some studies have been done on the comparison of these  
72 methods in simple BVPs like plane strain compression. It is found that the G&S gives better  
73 regularisation results than the GD one (Galavi and Schweiger, 2010; Guo and Stolle, 2013;  
74 Summersgill et al., 2017; Mallikarachchi and Soga, 2020; Gao et al., 2022). But the  
75 performance of these functions in real-world BVPs has not been evaluated.

76  
77 The main aim of this study is to carry out a comprehensive comparison of these functions in  
78 various BVPs, including drained and undrained plane strain compression, the response of strip  
79 footings on level ground and near a slope and a retaining wall (passive and active conditions).  
80 An anisotropic sand model accounting for the evolution of anisotropy is used. The increment  
81 of the void ratio which has a significant influence on the strain softening is assumed nonlocal.  
82 The three nonlocal weight functions and constitutive model will be first introduced. The  
83 performance of the regularisation methods will then be compared in different BVPs.

84

## 85 **2. The weight functions for nonlocal regularisation**

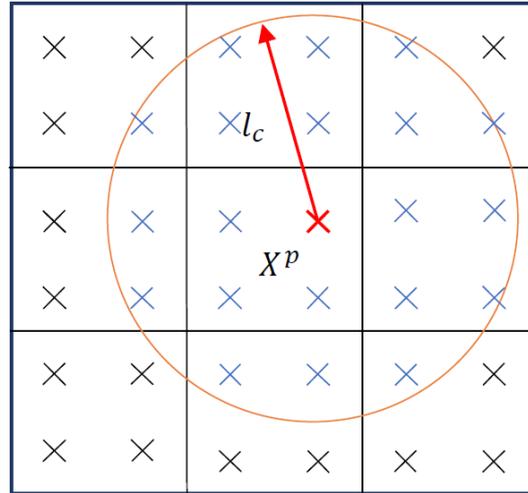
86 The GD function was first introduced by Eringen (1974) and then successfully implemented in  
87 damage models with strain softening (Bažant et al., 1984). It is expressed as

$$88 \quad \omega_i = \frac{1}{\sqrt{\pi}l_c} \exp\left(-\frac{r_i^2}{l_c^2}\right) \quad (1)$$

89 where  $\omega_i$  represent the weight function of integration point  $i$ ,  $r_i$  is the distance between the  
90 current integration point and the  $i - th$  integration point,  $l_c$  is a nonlocal parameter termed  
91 internal length which is dependent on the mean size of soil particles (Galavi and Schweiger,  
92 2010). Fig. 1 shows the physical significance of internal length in a 2D problem. Fig. 2 shows  
93 the plot of Eq. (1) in 1D condition. It is obvious that the GD function shows the highest  
94 contribution to the calculated nonlocal variable at the centre and diminishes along the

95 distance. As mentioned by Vermeer and Brinkgreve (1994), the nonlocal variable is  
 96 concentrated at the local point and cannot spread to surrounding points which has a negative  
 97 effect on the nonlocal method. This results in a centre concentration of the softening variable  
 98 when local strain is treated as a nonlocal variable.

99

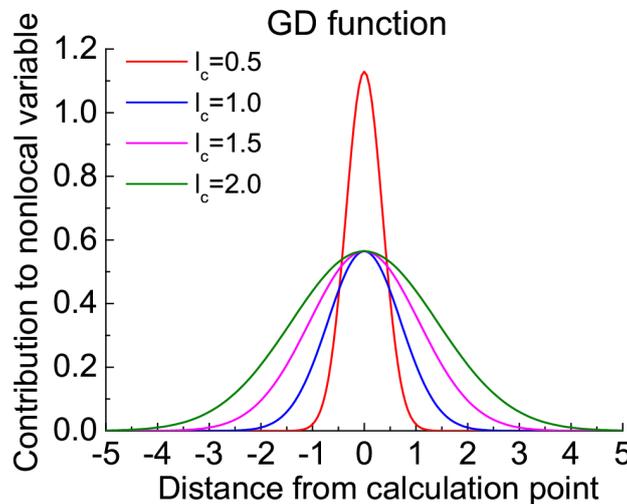


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**Fig. 1 Schematic diagram showing neighbouring integration point of  $X^p$**



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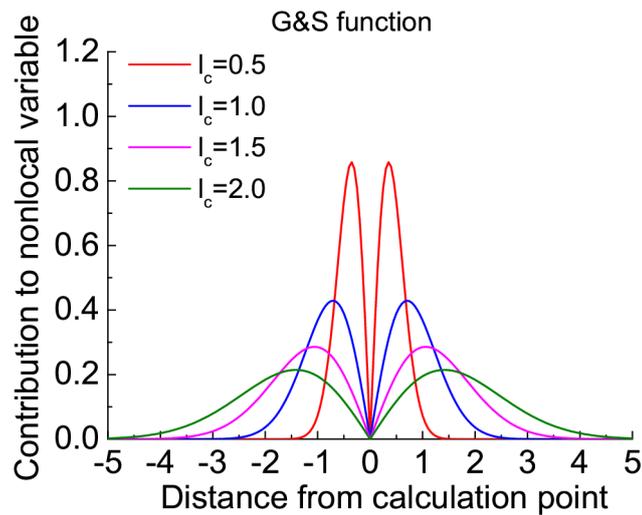
**Fig. 2 The Gaussian distribution function in 1D condition**

106 Based on the hypothesis that the deformation at a point is more influenced by the response  
 107 at the neighbourhood rather than the concentrated deformation at the point itself, Galavi  
 108 and Schweiger (2010) have proposed the following weight function

109

$$\omega_i = \frac{r_i}{l_c^2} \exp\left(-\frac{r_i^2}{l_c^2}\right) \quad (2)$$

110 As shown in Fig. 3, the contribution of the G&S weight function to the calculated nonlocal  
 111 variable is zero in the centre point and efficiently spreads from the concentrated local point  
 112 to a larger zone. This is different from the GD function with the maximum value at the centre.  
 113 In addition, the G&S weight function shows two same peaks with a distance of  $0.707l_c$  from  
 114 the centre (Galavi and Schweiger, 2010).  
 115



116  
 117

118 **Fig. 3 The Galavi and Schweiger (2010) distribution function in 1D condition**

119

120 Moreover, another method to overcome the limitations of Gaussian distribution is proposed  
 121 by Vermeer and Brinkgreve (1994), which is a linear combination of the local and the nonlocal  
 122 variables. A nonlocal parameter  $m$  is applied to change the nonlocal averaging formulation.  
 123 This method was called the over-nonlocal method. The nonlocal variable is expressed as

$$124 \quad \bar{\varpi}(x) = (1 - m)\varpi(x) + \frac{m}{V} \int_V \omega(x, \xi) \varpi(\xi) d\xi \quad (3)$$

125 where  $\bar{\varpi}(x)$  is the nonlocal variable and  $\varpi(x)$  is the local variable. The parameter  $m$  provides  
 126 the relative contribution from local and nonlocal parts. When  $m < 1$  in Eq. (3), the nonlocal  
 127 variable produces less effect than the local one. On the contrary, the contribution of the local  
 128 variable will be negative when  $m > 1$ . Existing research has shown that  $m > 1$  should be  
 129 used to achieve the best regularisation results (Vermeer & Brinkgreve, 1994; Lü et al., 2009;  
 130 Xue et al., 2022). But the exact value is dependent on the model and has to be determined  
 131 via trial and error.

132

### 133 3 The constitutive model and nonlocal regularisation

#### 134 3.1 Constitutive model

135 The constitutive model used here has been presented by Gao et al. (2022). Only the yield  
136 function and plastic hardening law which affect the strain softening are given here. The yield  
137 function is expressed as

$$138 \quad f = \frac{R}{g(\theta)} - H = 0 \quad (4)$$

139 where  $R = \sqrt{\frac{3}{2}r_{ij}r_{ij}}$ , with  $r_{ij} = (\sigma_{ij} - p\delta_{ij})/p$ ,  $\sigma_{ij}$  is the stress tensor,  $p = \sigma_{ii}/3$  is the  
140 mean effective stress,  $\delta_{ij}$  is the Kronecker delta ( $= 1$  for  $i = j$ , and  $= 0$  for  $i \neq j$ ),  $H$  is the  
141 hardening parameter and  $g(\theta)$  is an interpolation function which describes the variation of  
142 critical state stress ratio with the Lode angle  $\theta$  of  $r_{ij}$  (Li and Dafalias, 2002).

143

144 The hardening law for the yield function (evolution of for  $H$ ) is expressed as

$$145 \quad dH = \langle L \rangle r_H = \langle L \rangle \frac{Gh_1e^{h_2A}}{(1+e)^2\sqrt{pp_aR}} [M_c g(\theta)e^{-n\zeta} - R] \quad (5)$$

146 where  $h_1$ ,  $h_2$ , and  $n$  are three model parameters;  $A$  is the anisotropic variable (Li and Dafalias,  
147 2012);  $G$  is the elastic shear modulus;  $L$  is the loading index and  $\langle \ \rangle$  are the Macaulay  
148 brackets which make  $\langle L \rangle = L$  for  $L > 0$  and  $\langle L \rangle = 0$  for  $L \leq 0$ ;  $p_a$  is the atmospheric  
149 pressure (101 kPa);  $M_c$  is the critical state stress ration in triaxial compression;  $e$  is the void  
150 ratio, and  $\zeta$  is the dilatancy state parameter. The expression of  $g(\theta)$  can be found in Gao et  
151 al. (2022).

152

153 It is evident that several variables affect plastic hardening, and thus, strain softening (Gao et  
154 al., 2022). But only the increment of void ratio is assumed nonlocal for two main reasons (Gao  
155 et al., 2022). First, the void ratio is a key state variable that affects the behaviour of sand.  
156 Secondly, making the other state variables nonlocal can be computationally expensive. The  
157 increment of the nonlocal void ratio is expressed below

$$158 \quad de = (1 + e)d\varepsilon_{vn} \quad (6)$$

$$159 \quad d\varepsilon_{vn} = \frac{\sum_{k=1}^N w_i v_i d\varepsilon_{vi}}{\sum_{k=1}^N w_i v_i} \quad (7)$$

160 where positive  $de$  means volume contraction and  $d\varepsilon_{vn}$  is the nonlocal volumetric strain  
161 increment,  $N$  is the number of integration points within the averaging area,  $v_i$  and  $d\varepsilon_{vi}$

162 represent the volume and local volumetric strain increment of integration point  $i$ . Eq. (6) and  
 163 (7) can be used for the GD and G&S functions. When the over-nonlocal method is used, the  
 164 void ratio increment is expressed as

$$165 \quad de = (1 + e)[(1 - m)d\varepsilon_{vl} + md\varepsilon_{vn}] \quad (8)$$

166 where  $d\varepsilon_{vl}$  is the total local volumetric strain increment for each step. In Eq. (8),  $d\varepsilon_{vn}$  is  
 167 calculated using the GD function.

168

### 169 **3.2 Implementation of the nonlocal method**

170 The model used the explicit stress integration method (Zhao et al., 2005; Gao and Zhao, 2013;  
 171 Zhou, et al., 2021; Zhou, et al., 2022; Lu, et al. 2023) and two user subroutines, UMAT (user-  
 172 defined materials) and USDFLD (user-defined field variables), are needed for implementing  
 173 the nonlocal method in Abaqus.

174 To increase the computation efficiency, a scaling variable  $r_v$  defined as below is introduced  
 175 (Gao et al., 2021)

$$176 \quad r_v = \frac{d\varepsilon_{vn}}{d\varepsilon_{vl}} \quad (9)$$

177 where  $d\varepsilon_{vn}$  is the nonlocal volumetric strain increment and  $d\varepsilon_{vl}$  is the total local volumetric  
 178 strain increment for each increment. Both of them are calculated before the sub-increment  
 179 of strain is applied in the stress integration.

180

181 At the end of each sub-increment the void ratio increment  $de^s$  is

$$182 \quad de^s = (1 + e)d\varepsilon_{vl}^s r_v \quad (10)$$

183 where  $d\varepsilon_{vl}^s$  is the local sub-increment of volumetric strain (Gao et al., 2022). The summation  
 184 of  $de^s$  at all sub-increments is the total void ratio change in a step. For the over-nonlocal  
 185 method,  $de^s$  is calculated below

$$186 \quad de^s = (1 + e)[(1 - m)d\varepsilon_{vl}^s + md\varepsilon_{vl}^s r_v] \quad (11)$$

187

### 188 **4. Plane strain compression tests**

189 The sample used in this study is 60 mm wide and 120 mm high as shown in Fig. 4. The  
 190 boundary condition is shown in Fig. 4. A confining pressure of  $p_0 = 200\text{kPa}$  is applied on the  
 191 two vertical sides. Vertical displacement is applied on the top side with the horizontal  
 192 displacement unconstrained. The bottom side is pinned at the left and free to move to the

193 right. A square ‘weak’ area (12mm×12mm) with inclined bedding plane orientation ( $\alpha = 45^\circ$ )  
194 is implemented, which is used to trigger a shear band in the plane strain compression test.  
195 For the remaining part of this specimen, the bedding plane orientation is horizontal and  $\alpha =$   
196  $0^\circ$ . The anisotropic model parameters are shown in Table 1. All the parameters are the same  
197 as those in Gao et al. (2022). Note that the parameter  $m_d$  is  $m$  in Gao et al. (2022). The initial  
198 void ratio of the sample is  $e_0 = 0.65$  (relative density  $D_r = 85.6\%$ ), and the initial degree of  
199 anisotropy is  $F_0 = 0.4$ . All simulations in this study have used 8-noded plane strain quadratic  
200 elements with reduced integration (CPE8R). Note that all the simulations to be presented  
201 below use this element. The thickness of the soil is assumed 1m in processing the results.

202

203 The internal length  $l_c$  is an important parameter for nonlocal regularisation models. The size  
204 of the internal length determines how many integration points can be involved in nonlocal  
205 regularisation.  $l_c$  should be equal to or larger than the maximum mesh size to make sure that  
206 sufficient integration points are involved. Bigger  $l_c$  means that the stress and strain  
207 relationship of the current integration point is affected by that of integration points further  
208 away. Fig. 5 illustrates the effect of  $l_c$  on the vertical reaction force and displacement curves  
209 simulated by the different weight functions. In these models, the mesh size of 0.004 m was  
210 selected under drained conditions. The  $l_c$  does not affect the solutions before the peak  
211 reaction force. Higher peak vertical reaction force and a slower rate of strain-softening were  
212 obtained by increasing  $l_c$  during post-peak. Furthermore, the GD and ON functions predict a  
213 slower rate of the strain-softening curve than the GD function. The internal length determines  
214 the range within which the integration points are considered in the nonlocal averaging. When  
215 it is bigger, more integration points are accounted for in the weight functions of each  
216 integration point. This means that the local load is artificially distributed to more neighbouring  
217 integration points, leading to a lower rate of strain softening. In the simulations for plane  
218 strain compression below,  $l_c = 0.012 m$  is used. It should be mentioned that  $l_c$  also has  
219 influence on the shear band thickness, which will be shown in subsequent sections. The real  
220 shear band thickness of sand is about  $10-20d_{50}$ , where  $d_{50}$  is mean particle size (Galavi and  
221 Schweiger, 2010). If the real shear band thickness were to be matched in FE modelling, very  
222 small mesh size has to be used because the shear band thickness is close to  $l_c$ . This would  
223 cause issues like excessive computation time and numerical divergence. Therefore, proper  $l_c$

224 is typically chosen based on the size of solution domain, which can guarantee mesh-  
225 independent results but not realistic shear band thickness.

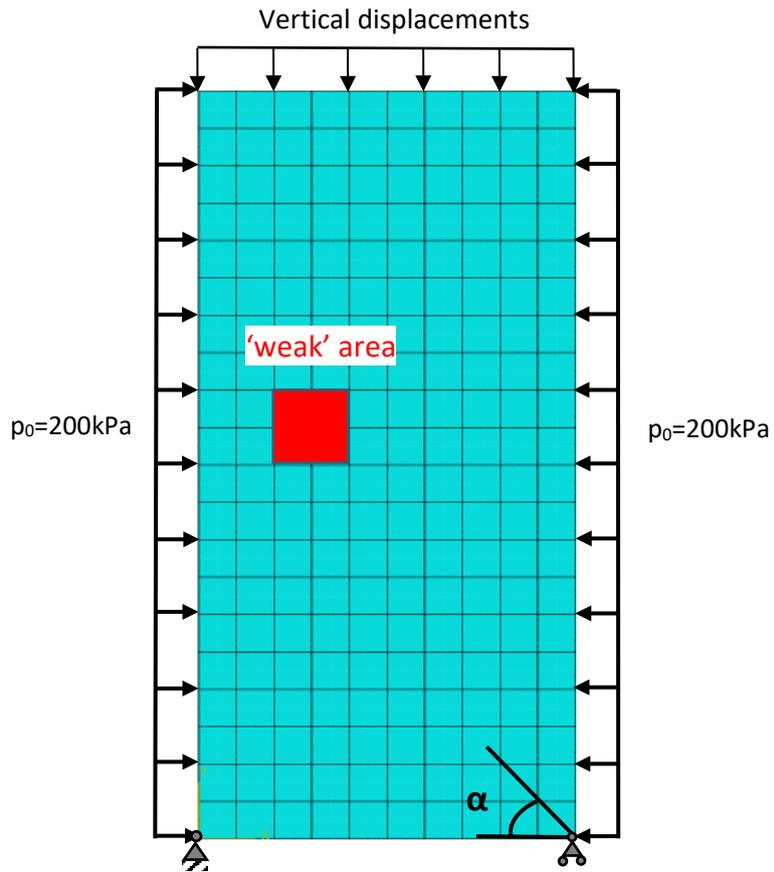
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**Table 1 Model parameters for Toyoura sand**

Parameters	Value
$G_0$	125
$\nu$	0.1
$M_c$	1.25
$c$	0.75
$e_\Gamma$	0.934
$\lambda_c$	0.019
$\xi$	0.7
$n$	2.0
$h_1$	0.45
$d_1$	1.0
$m$	3.5
$k_f$	0.5
$e_A$	0.075
$k_h$	0.03
$h_2$	0.5

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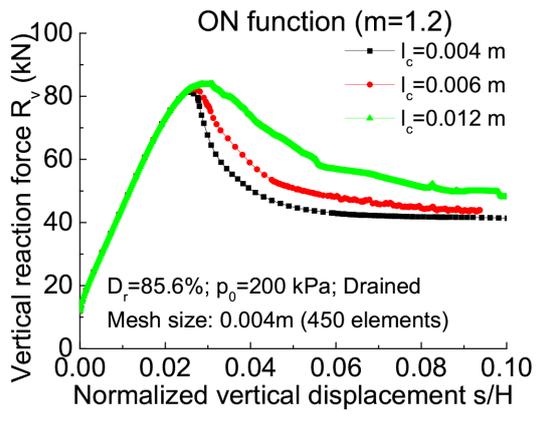
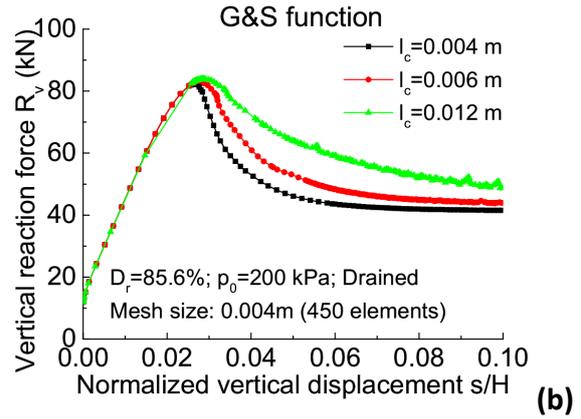
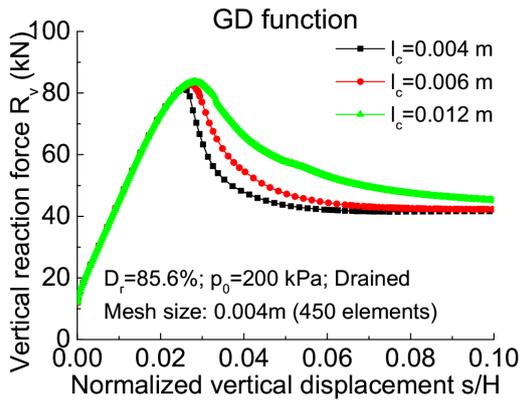


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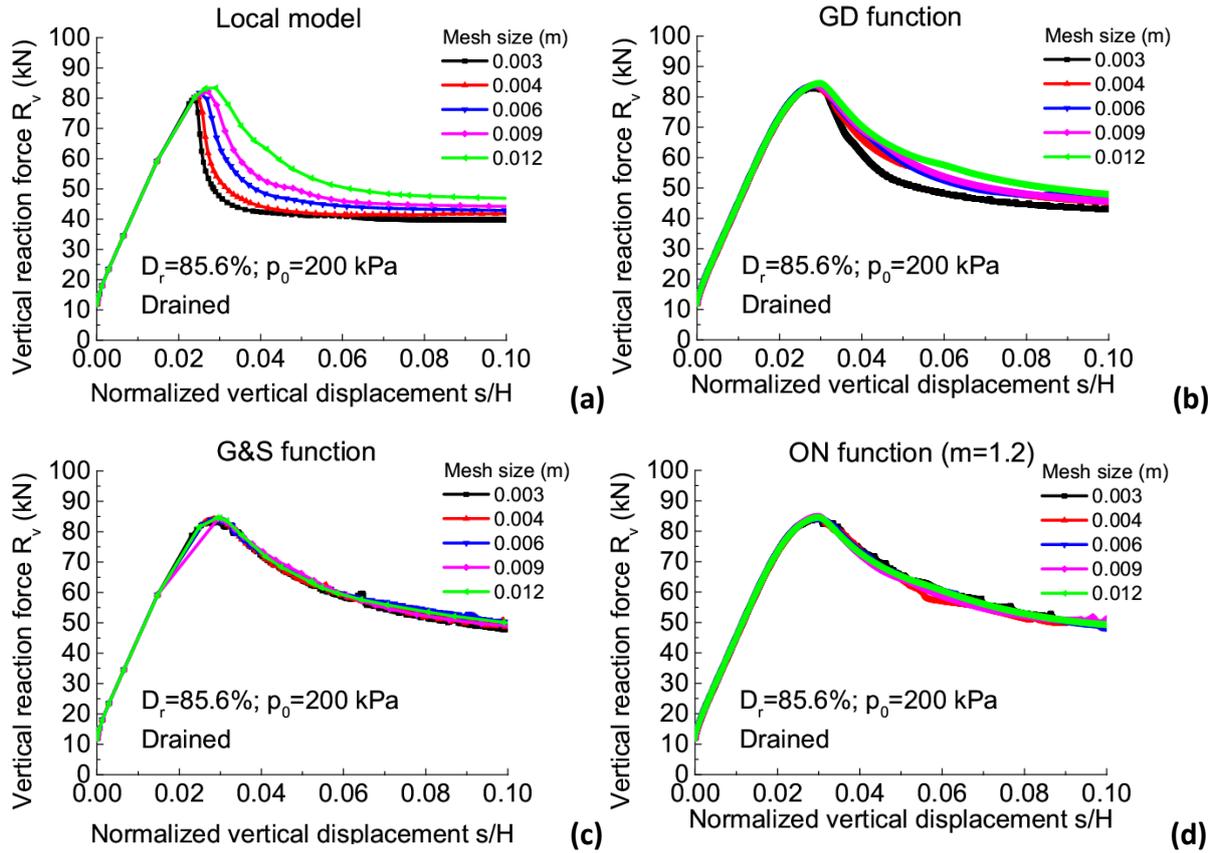
**Fig. 4 The boundary conditions and bedding plan orientation for the plane strain compression simulations**



**Fig. 5 The effect of internal length on the force-displacement relationship in drained plane strain compression test: (a) GD function; (b) G&S function; (c) ON function**

#### 4.1 Drained plane strain compression tests

Fig. 6 shows the force-displacement curves predicted by the local model and three nonlocal models. The GS and ON functions give better regularisation results than the GD one. The main reason is that the local variable has more significant influence on the results when the GD function is used. It should be mentioned that  $m = 1.2$  is chosen for the ON method through trial-and-error. Smaller  $m$  gives mesh-dependent solutions, but higher  $m$  causes numerical divergence in the simulations. In the strip footing problem to be discussed in the subsequent sections, higher  $m$  is found to give steep reduction of reaction force acting on the footing after peak, which is not consistent with the experimental observations in centrifuge tests.

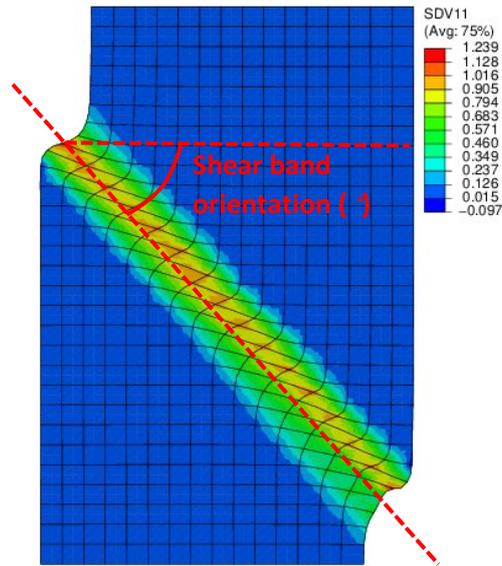


**Fig. 6 Comparison of the local and nonlocal models on the force-displacement relationship for drained plane strain compression: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**

The orientation of the shear band ( $\beta$ ) is directly measured from shear strain contours as shown in Fig. 7. The predicted angle of shear band orientation decreases as the mesh size increases (Fig. 8). All the nonlocal functions reduce but cannot eliminate the mesh dependency of shear band orientation. This could be due to that only one variable that affects the strain softening is assumed nonlocal. The mesh dependency could be further reduced if more state variables in the hardening law are assumed nonlocal. But this would significantly reduce the computation efficiency.

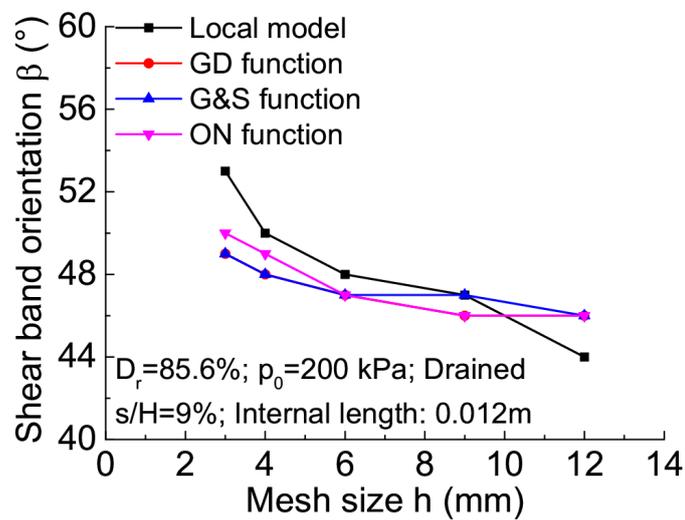
The thickness of shear band is measured based on the shear strain distribution across a shear band at  $s/H = 7\%$  (Fig. 9). Fig. 10 shows the determination of shear band thickness for the nonlocal model. In Fig. 10 (a)  $t_{s1}$  and  $t_{s2}$  represent shear band thickness for mesh size 0.004 m and 0.006 m, respectively. The effect of mesh size on shear band thickness is shown in Fig. 11 (a). The shear band thickness simulated by the local model increases significantly with the

265 mesh size. The nonlocal models give a small variation of shear band thickness when the mesh  
 266 size  $h < l_c$ . All nonlocal models give the same shear band thickness as that of the local model  
 267 when size  $h = l_c$ . The shear band thickness predicted by the nonlocal models increases with  
 268  $l_c$  (Fig. 11b) and the ON model predicts wider shear bands.  
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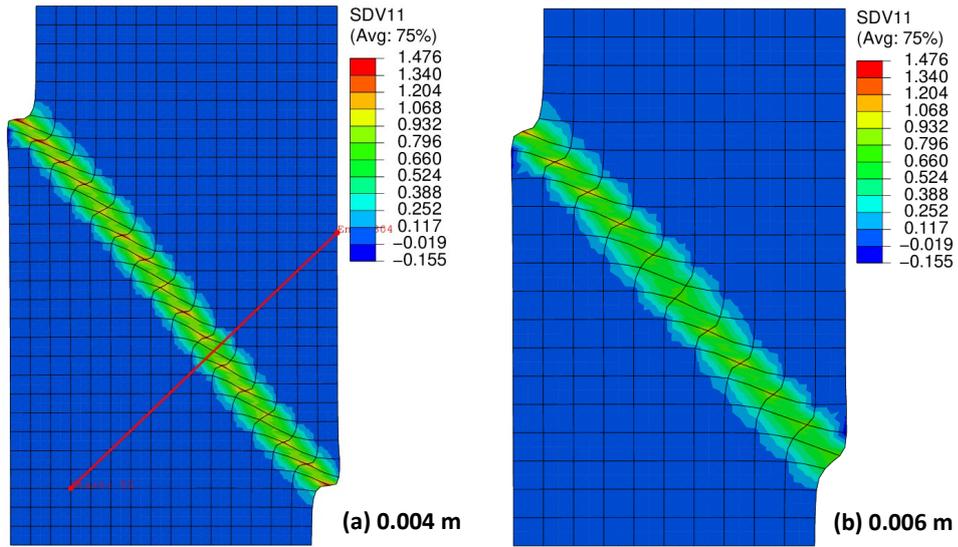
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**Fig. 7 Shear strain contour for measuring the shear band orientation**



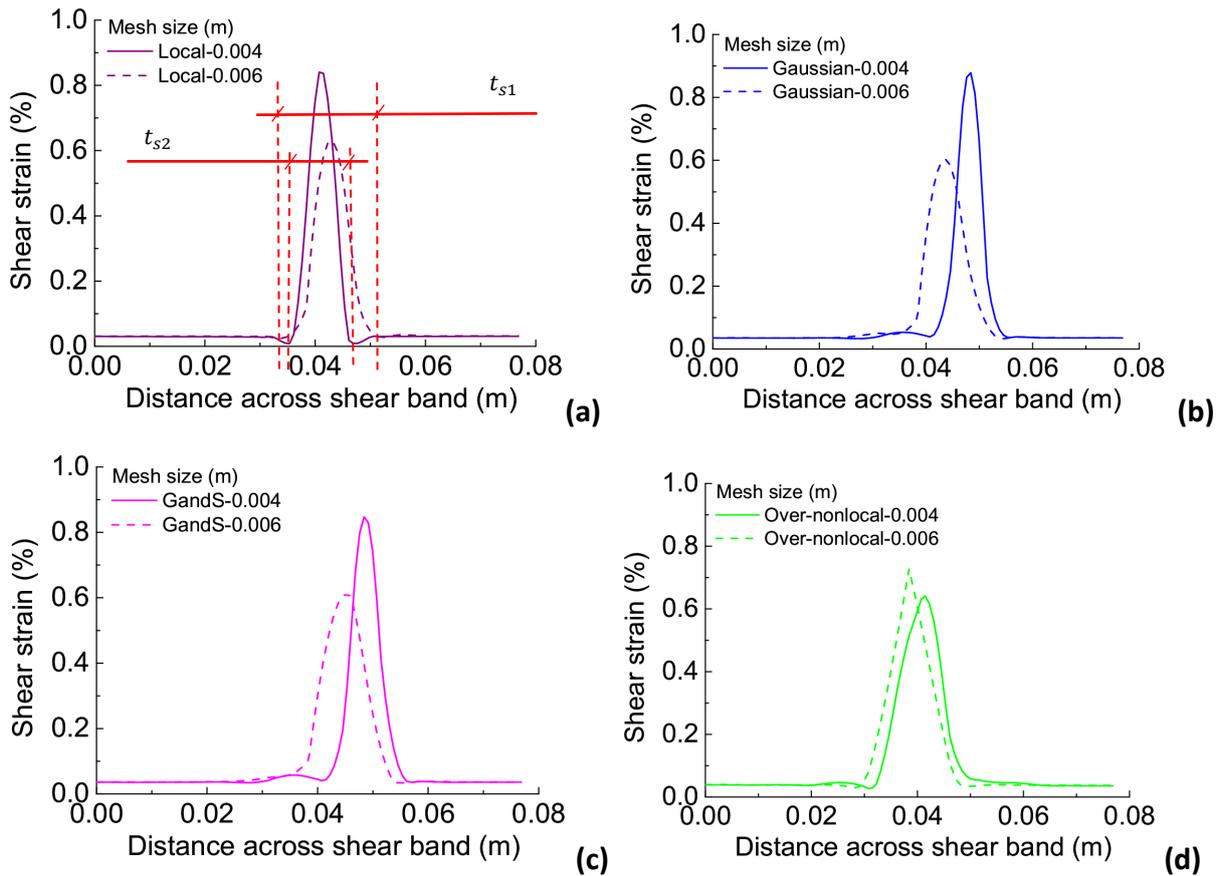
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**Fig. 8 Comparison of shear band orientation for drained plane strain compression test**



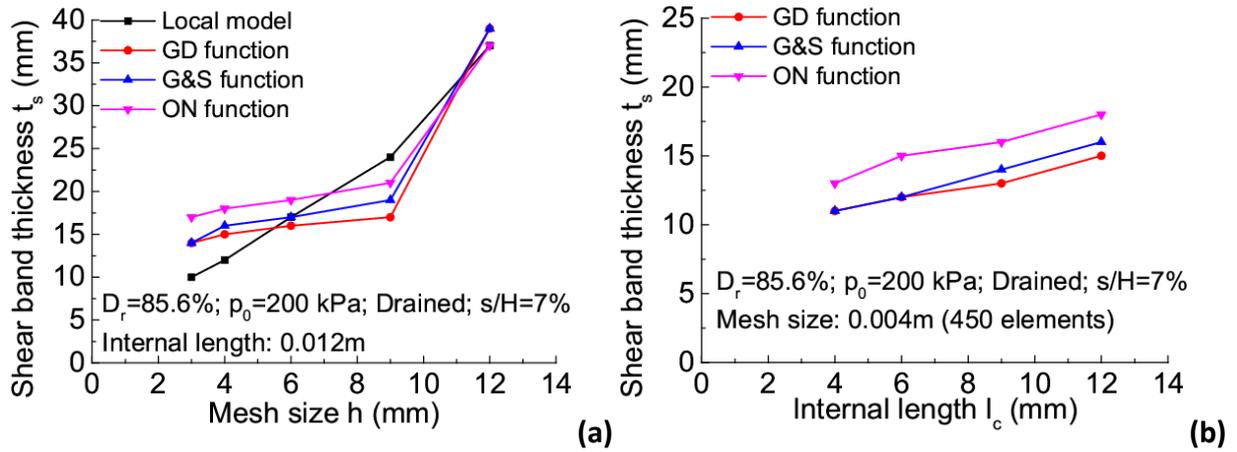
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**Fig. 9 Cross-section contour based on the shear strain under drained condition with mesh size of (a) 0.004 m and (b) 0.006 m for the local model**



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**Fig. 10 Comparison of cross-section profiles based on the shear strain: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**



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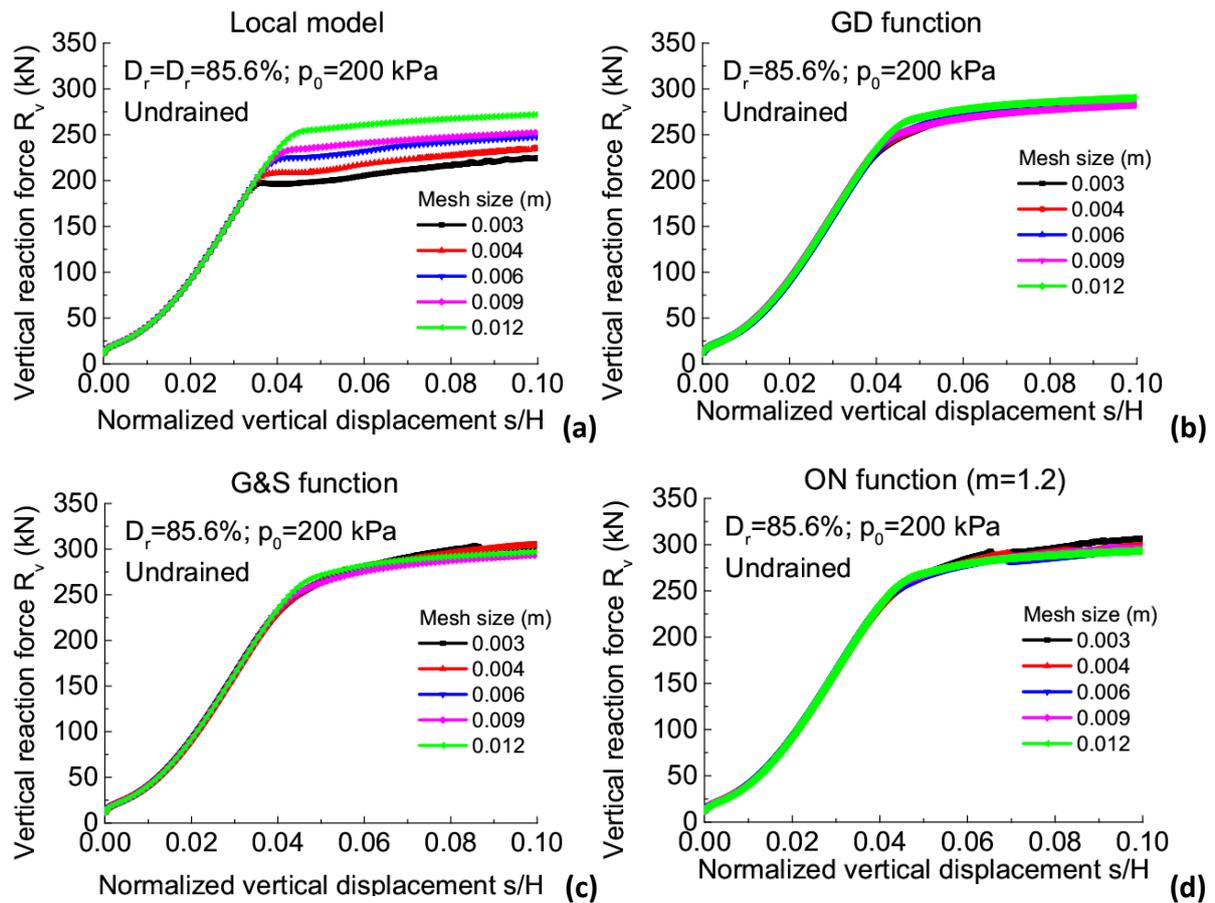
284 **Fig. 11 Comparison of the effect of (a) mesh size and (b) internal length on the shear band**  
 285 **thickness in plane strain compression**

286

## 287 4.2 Undrained plane strain compression tests

288 In undrained plane strain compression, the permeability of soil is set very small and water  
 289 flow at all boundaries is closed. Fig. 12 shows the relationship between vertical displacement  
 290 and reaction force for different models. It is evident that the nonlocal models give mesh-  
 291 independent results.

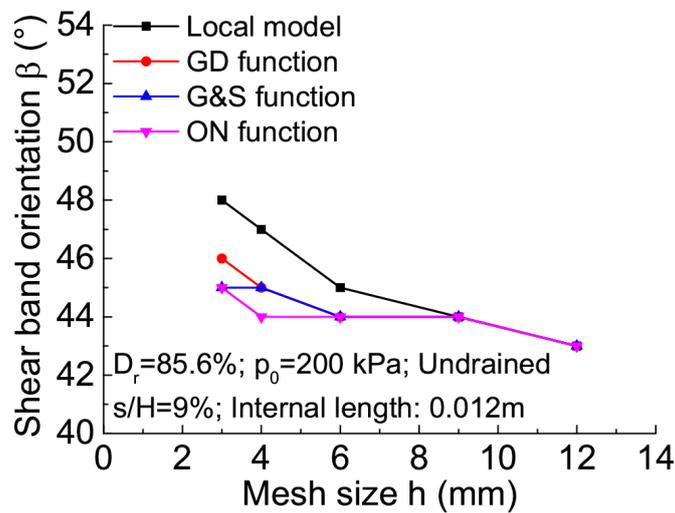
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**Fig. 12 Comparison of the force-displacement relationship for undrained plane strain compression test: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**

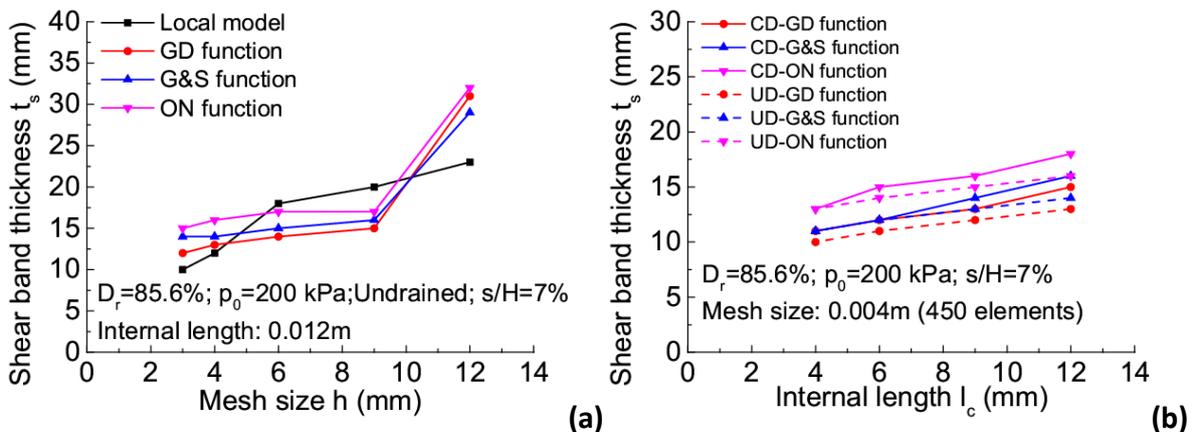
The shear band orientation in undrained plane strain compression increases when the mesh is refined for all models. The mesh dependency can be reduced but not eliminated by the nonlocal treatment (Fig. 13). It is worth noting that the nonlocal models give the same shear band orientation as the local model when the mesh size is greater than 0.009 m. The nonlocal models also give a bigger variation of shear band orientation in undrained tests than in drained tests. The main reason is that there is a smaller change in the void ratio in an undrained test, which makes the nonlocal regularisation using the void ratio less effective. The shear band thickness predicted by the models is shown in Fig. 14. Similar to the drained cases, the nonlocal models give a small variation of shear band thickness when the mesh size is smaller than the internal length. But the shear band thickness predicted by the nonlocal models at  $h = l_c$  is bigger than that of the local model. Moreover, it is found that the drainage condition has little influence on the shear band thickness at different internal lengths (Fig. 14b).

311 Fig. 15 shows the force and displacement relationship for loose sand in undrained plane strain  
 312 compression. Though the vertical reaction force decreases after the peak, the mesh size has  
 313 little influence on the results when the original model is used (Fig. 15a). The nonlocal models  
 314 give similar results (Fig. 15c-d). The reason is that the stress ratio of soil elements keeps  
 315 increasing though the deviator stress decreases. This is a strain-hardening response based on  
 316 the model, as increasing stress ratio means increasing hardening parameter  $H$  (Eq. 5). In  
 317 coupled dynamic loading (e.g., earthquake), the soil response will be a combination of that in  
 318 Fig. 12 and Fig. 15, wherein the nonlocal regularisation method is found to work. Therefore,  
 319 it is expected that the nonlocal regularisation technique also works for coupled analysis in  
 320 earthquakes.  
 321



322  
 323 **Fig. 13 Shear band orientation predicted by different models in undrained plane strain**

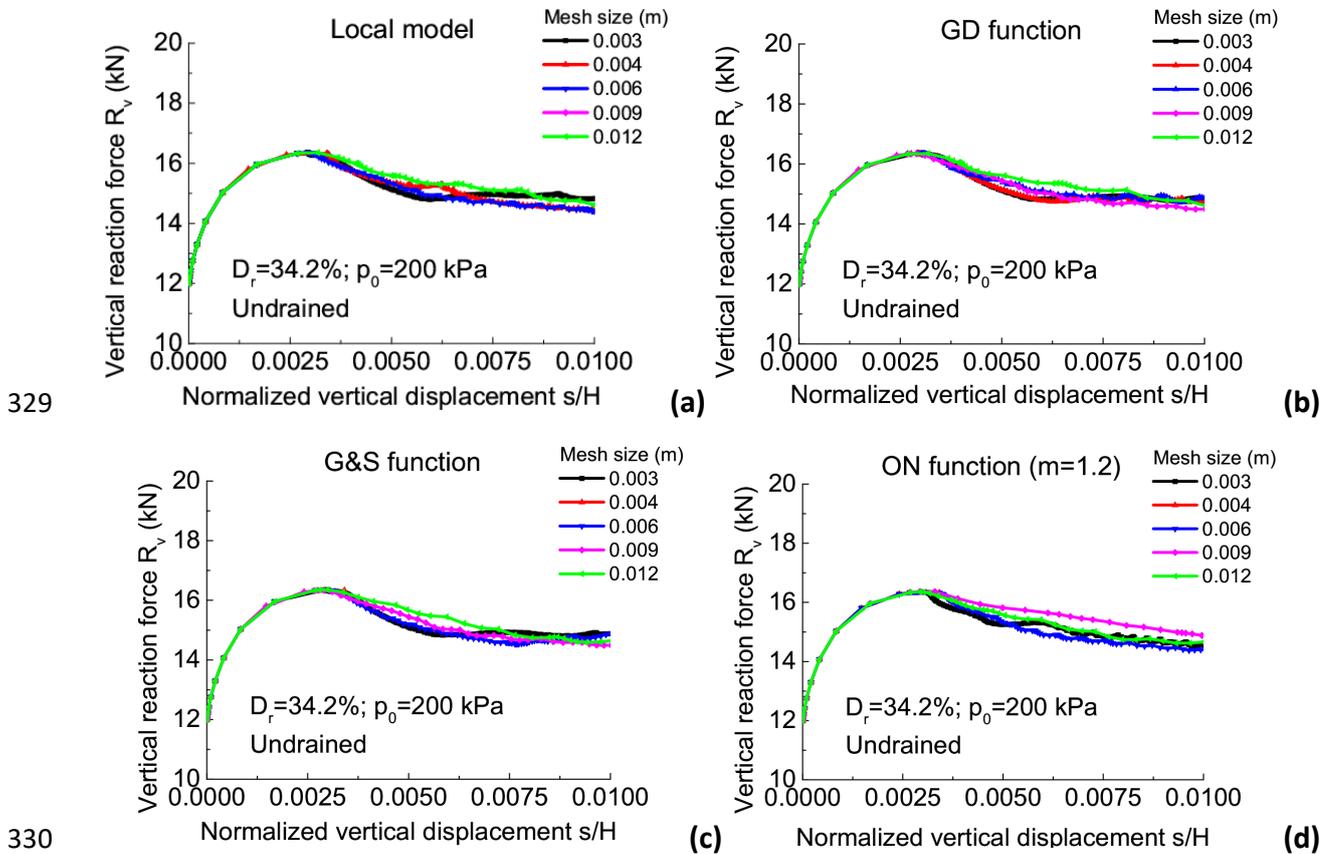
324



325

326 **Fig. 14 Comparison of the effect of (a) mesh size and (b) internal length on the shear band**  
 327 **thickness under undrained condition**

328



331 **Fig. 15 Comparison of the force-displacement relationship for static limited liquefaction**  
 332 **test: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**

333

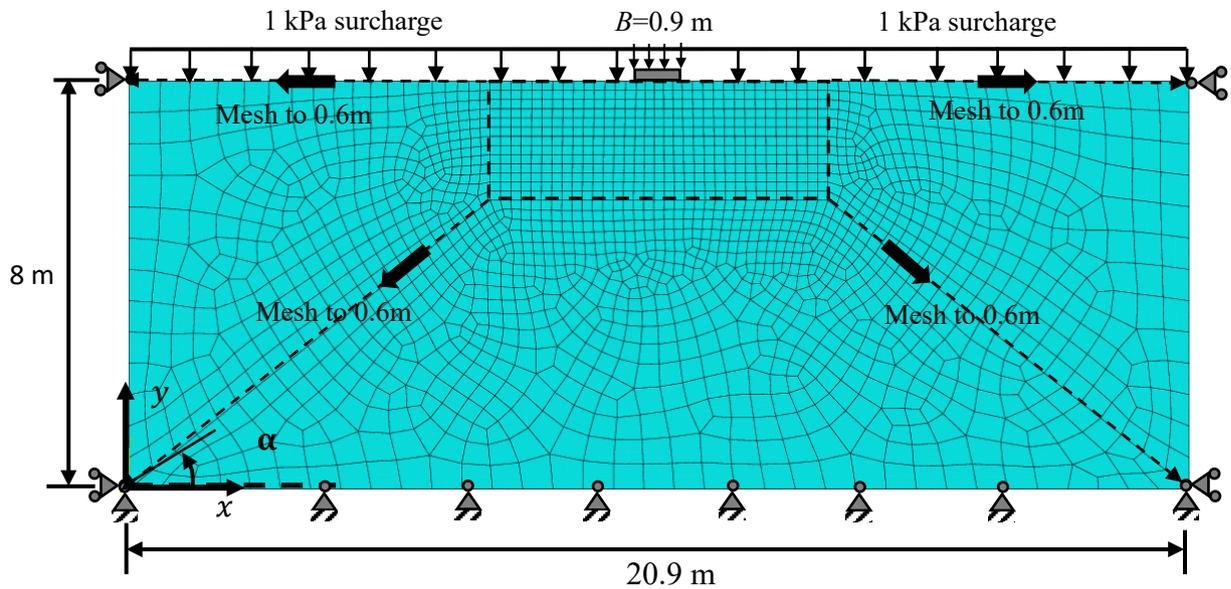
### 334 5. Strip footing problem

#### 335 5.1 Strip footing on level sand ground

336 The dimension of the strip footing problem is shown in Fig. 16. The footing with  $B=0.9$  m is  
 337 deformed by applying a uniform vertical deformation. The horizontal displacement is fixed to  
 338 simulate rough footings. Constant vertical pressure (1 kPa) is applied on the top surface to  
 339 avoid soil collapse with zero mean effective stress. The initial lateral earth pressure coefficient  
 340  $K_0 = 0.4$  (Okochi and Tatsuoka 1984), and the effective weight of Toyoura sand is  $\gamma' =$   
 341  $16\text{kN/m}^3$  as there is no water in the sand. Two sides of the sample are horizontally fixed,  
 342 while both horizontal and vertical movement is restricted for the bottom boundary. Details  
 343 can be found in Gao et al. (2020). Since the vertical load and vertical settlement relationship  
 344 is mainly affected by the rectangle area beneath the footing, hence, the mesh size far away

345 from the footing is setting a fixed value (0.6m) for all models. The bedding plane orientation  
 346 is horizontal and  $\alpha = 0^\circ$ . The relative density  $D_r = 85.6\%$  and the initial degree of anisotropy  
 347 is  $F_0 = 0.4$ .

348



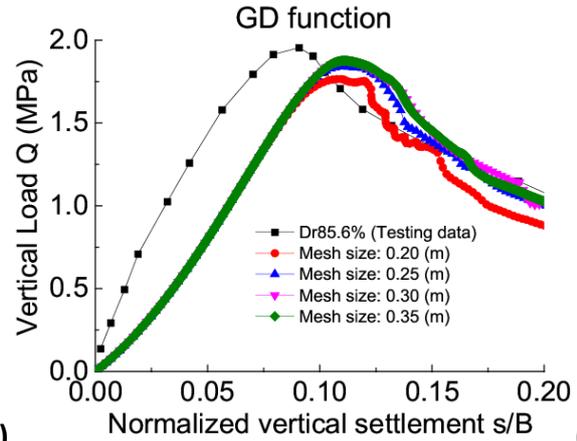
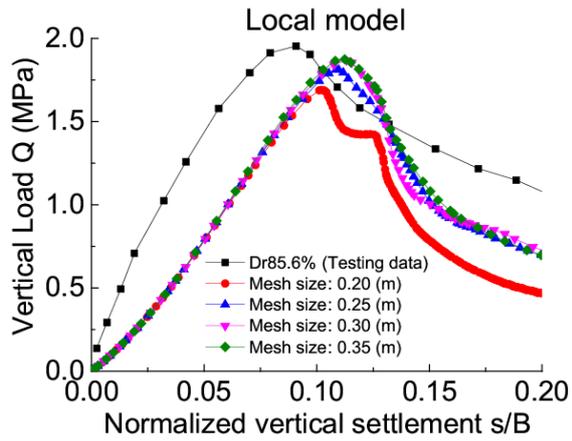
349

350 **Fig. 16 The boundary conditions of strip footing problem**

351

352 Fig. 17 shows the prediction of local and nonlocal models with  $l_c = 0.8 m$ . It is evident that  
 353 the local model prediction is highly mesh dependent. The G&S function gives the least mesh-  
 354 dependent results. The strain softening predicted by the GD and G&S functions are in good  
 355 agreement with the centrifuge test data (Kimura et al., 1985). But the ON model give very  
 356 steep reduction of  $Q$  after the peak, which does not match the experimental observations.  
 357 There are two reasons for this. First, this method gives excessive volume expansion of sand  
 358 under the strip footing (Fig. 18). Location of the elements in Fig. 18 is shown in Fig. 19. The  
 359 GD and GS models give similar prediction of void ratio evolution, while the void ratio increase  
 360 predicted by the ON model is about 90% higher. Higher void ratio causes lower strength and  
 361 failure of some elements, which lead to fast reduction of  $Q$ . Secondly, the ON method  
 362 assumes that the local variable makes negative contribution to the local one, which may not  
 363 be realistic. In this cases, such assumption causes failure or lower shear strength of more sand  
 364 elements. Moreover, it is found that bigger  $m$  value gives even steeper strain softening curve  
 365 for the ON function. Therefore, the ON function should not be used for this problem.

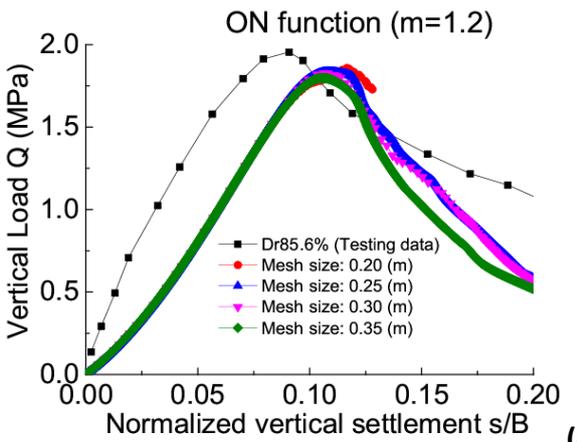
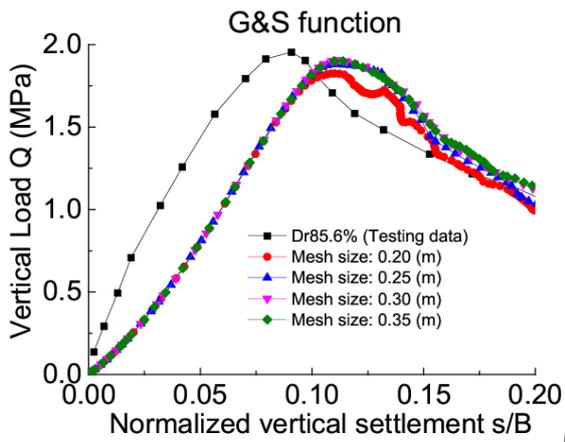
366



367

(a)

(b)



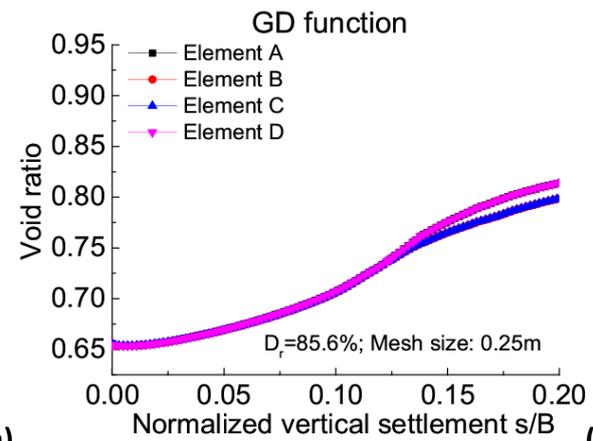
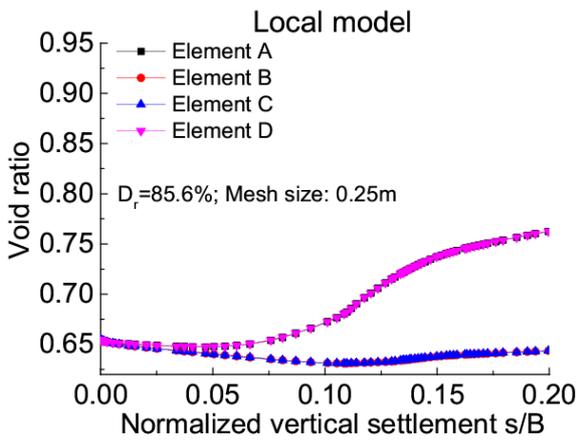
368

(c)

(d)

369 **Fig. 17 The comparison of the strip footing response on the sand with horizontal bedding**  
 370 **plane: (a)Local model; (b) GD function; (c) G&S function; and (d) ON function**

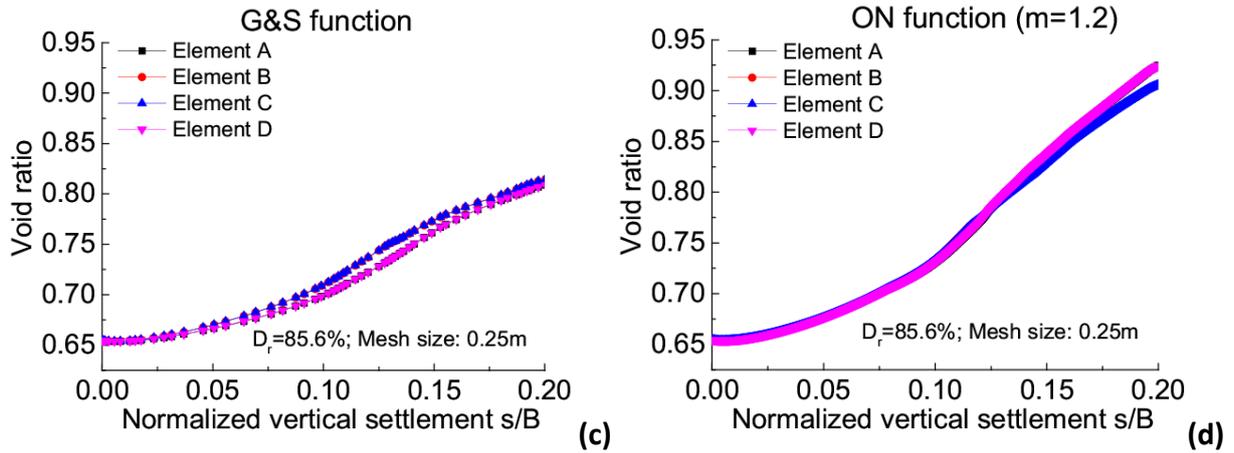
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372

(a)

(b)



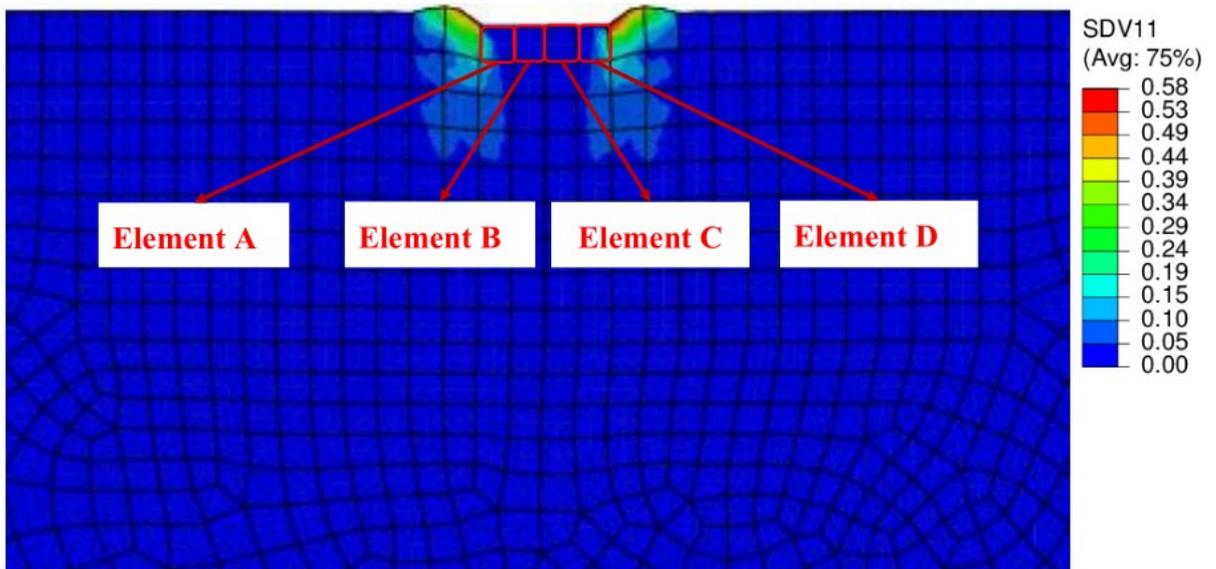
373

374

**Fig. 18 Comparison of void ratio evolution for elements under the strip footing: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**

375

376



377

378

**Fig. 19 Location of elements under the strip footing**

379

## 380 5.2 Strip footing near a sand slope

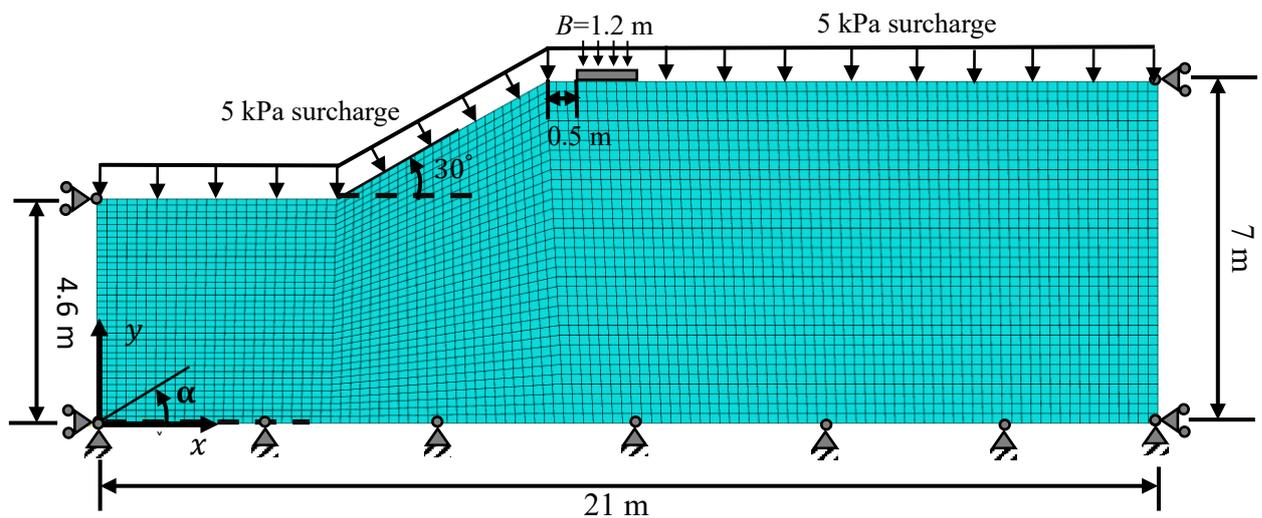
381 This problem is based on the simulations in Gao et al. (2021). The slope dimension and  
 382 boundary conditions are shown in Fig. 20. Since the ground surface is not level, a constant  $K_0$   
 383 cannot be applied. Therefore, the gravity loading method is used to generate the initial stress  
 384 state (Gao et al., 2021). First, gravity is applied on the same soil body by assuming that the  
 385 soil is elastic with a Poisson's ratio of  $\nu = 0.286$ , making  $K_0 = 0.4$  for a flat ground surface  
 386 (Gao et al., 2021). After that, the stress state is extracted and imported to the model as the

387 initial stress, which is used for the subsequent modelling. The soil density, initial void ratio  
388 and degree of anisotropy are the same as those in Fig. 16.

389

390 Fig. 21 shows the prediction of local and nonlocal models with  $l_c = 0.4 m$ . The local model  
391 gives different peak bearing capacity and  $s - Q$  curves after the peak as the mesh size  
392 changes (Fig. 21a). Nonlocal regularization reduces the mesh-dependent of  $s - Q$  curves (Fig.  
393 21c-d). The rate of strain softening is also reduced due to the nonlocal averaging of void ratio  
394 increment. To further reduce the mesh sensitivity, more nonlocal variables could be used, but  
395 this may increase the complexity of the model formulations and its implementation. Fig. 22  
396 shows the contour of shear strain distribution in the soil after the state for G&S function. A  
397 clear slip surface can be seen, which is independent of the mesh size.

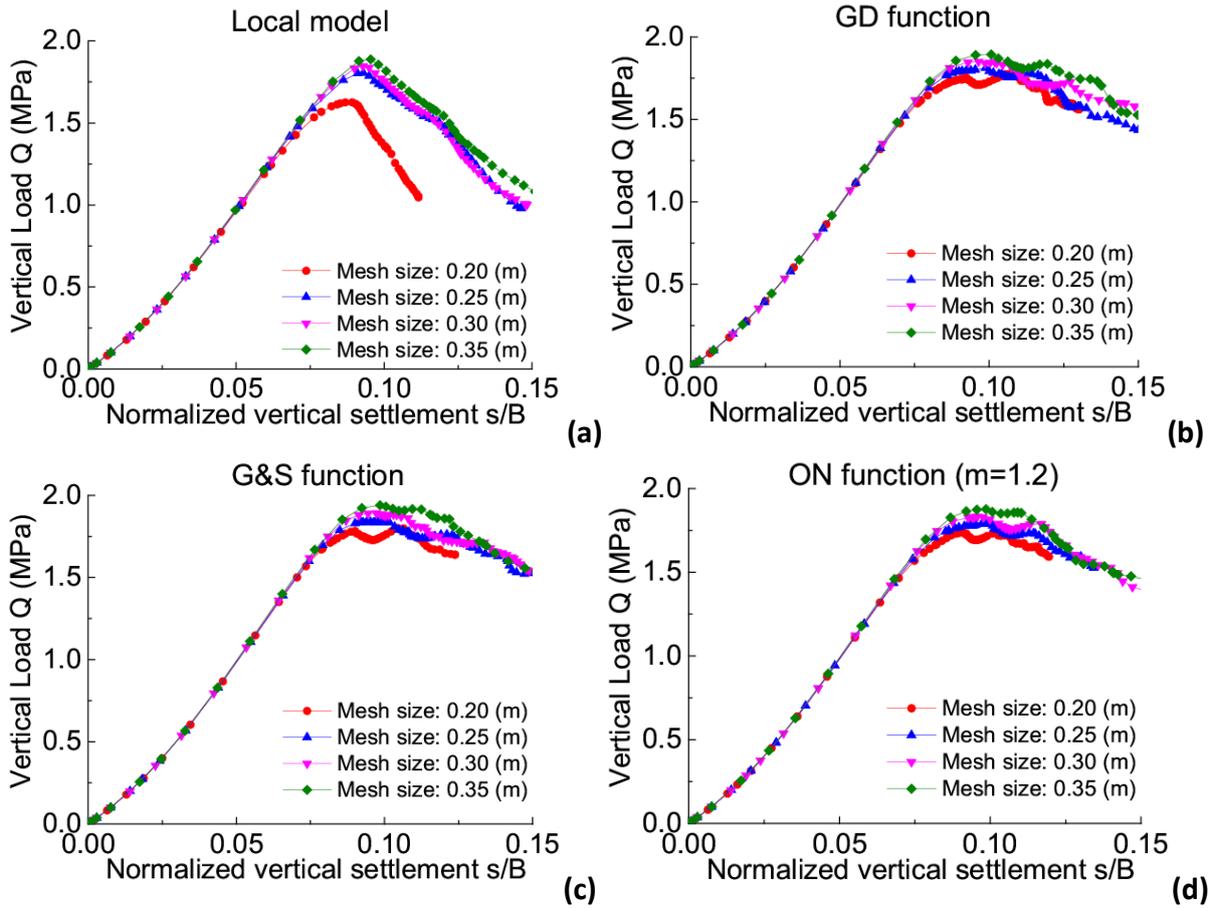
398



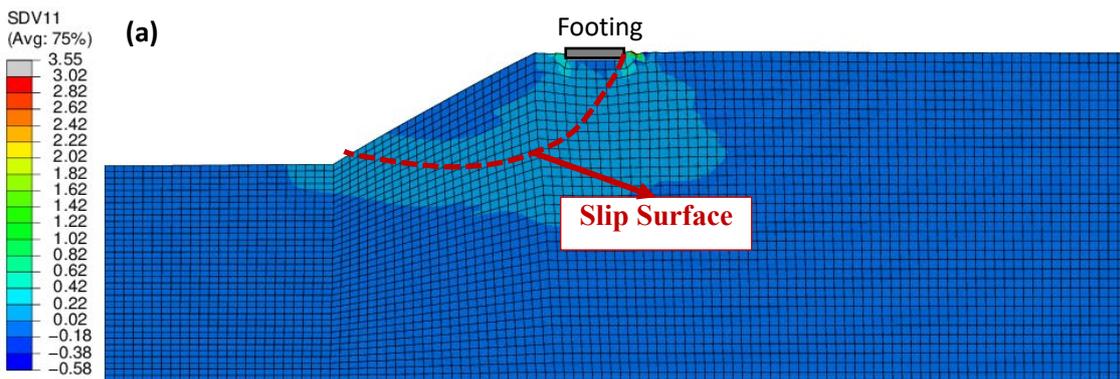
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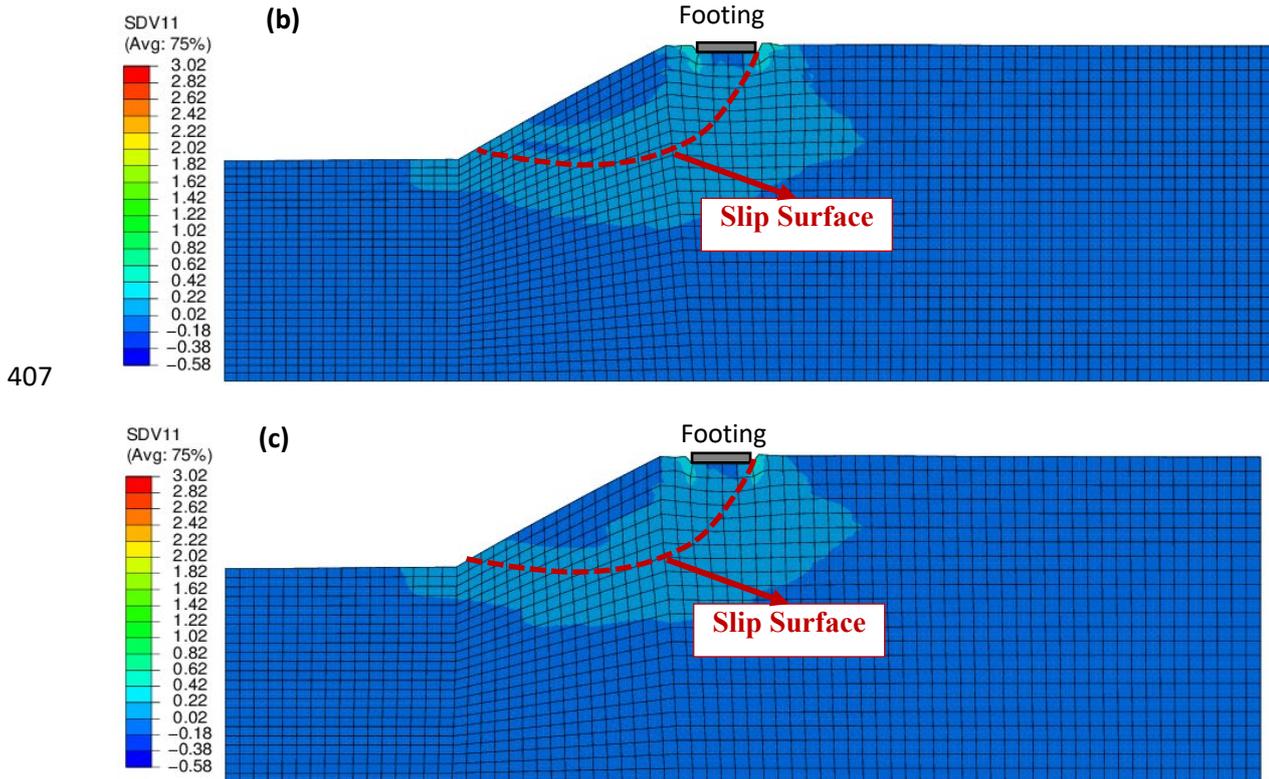
400

**Fig. 20 The boundary conditions of the strip footing near a slope**



401  
402  
403 **Fig. 21 The comparison of the strip footing response near a slope with horizontal bedding**  
404 **plane: (a) Local model; (b) GD function; (c) G&S function; (d) ON function**  
405





**Fig. 22** Shear strain distribution in the soil predicted by the G&S model at at  $s/B = 0.12$  with different mesh sizes: (a) 0.20 m; (b) 0.25 m; (c) 0.30 m

## 6. Response of retaining wall for level sand ground

412

413 Fig. 23 shows a soil domain measuring 10 m in length and 4.5 m in depth, with a rigid retaining

414 wall positioned on the right side of the backfill soil. The wall has a height of  $h_w = 4$  m and is

415 assumed to have an ideally smooth surface that prevents the transmission of shear stresses

416 at the interface with the soil. The retaining wall can undergo passive and active horizontal

417 translation, with passive movement towards the backfill and active movement away from it.

418 The bottom, left-side, and right-side boundaries are fully fixed. In all simulations, the bedding

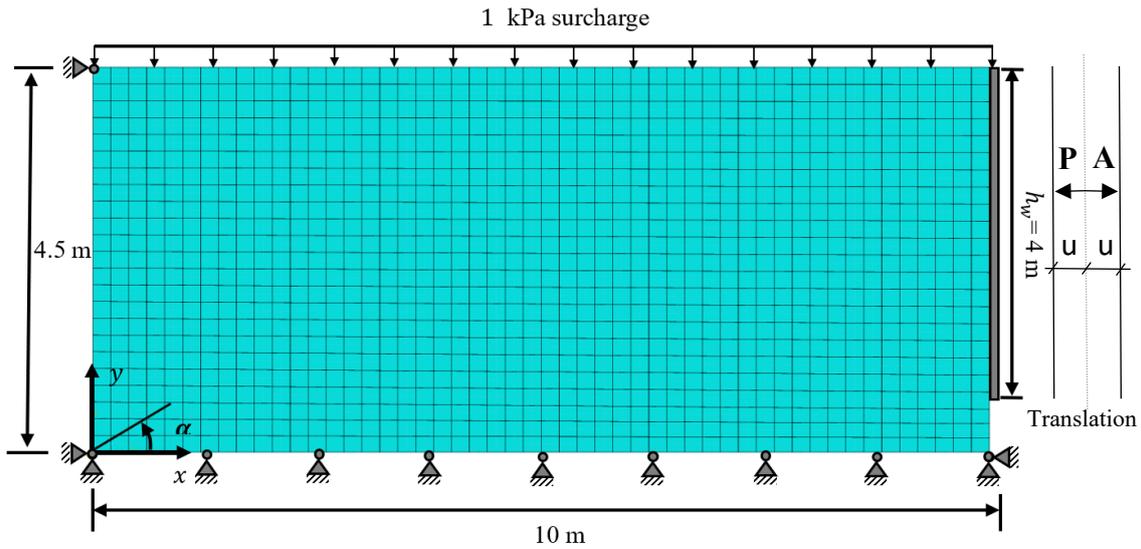
419 plane orientation is horizontal ( $\alpha = 0^\circ$ ) and the gravity is applied to the backfill soil while the

420 top surface of the backfill soil is subjected to a uniformly distributed surcharge of 1 kPa. The

421 same soil conditions as in Fig. 16 are used.  $l_c = 0.8$  m is applied for all simulations. The lateral

422 earth pressure is expressed as  $\sigma_h$  and the wall displacement  $u$  is normalised by the height of

423 the wall  $h_w$ .



424

425

**Fig. 23 The boundary conditions of the retaining wall problem**

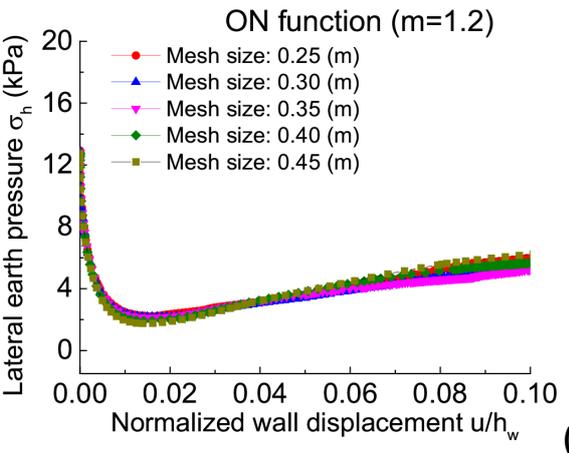
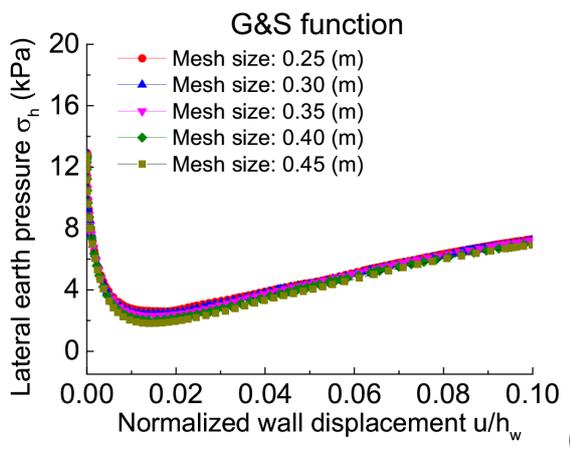
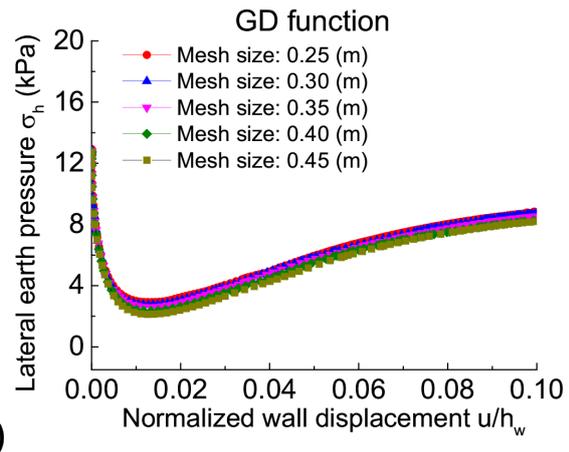
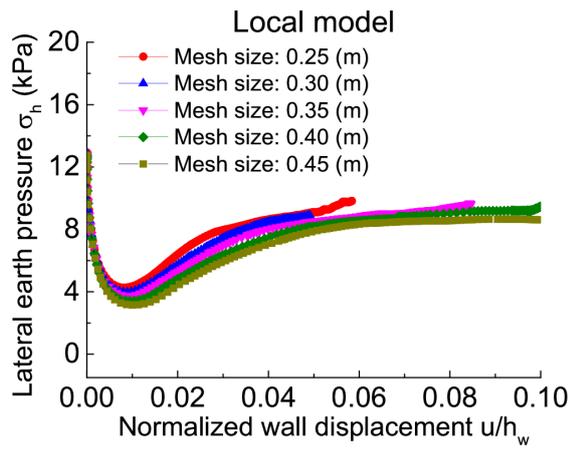
426

427 Fig. 24 shows the evolution of  $\sigma_h$  for the active condition. The local model gives mesh-  
 428 dependent  $\sigma_h - u/h_w$  curves. The G&S and ON functions are more efficient than GD function  
 429 in reducing the mesh-dependency.  $\sigma_h$  reaches the smallest value at  $u/h_w \approx 0.015$  and then  
 430 increases with  $u/h_w$ . This is caused by the strain softening of sand. Similar results have been  
 431 reported in Nübel and Huang (2004), Widulinski et al. (2011) and Guo and Zhao (2015). For  
 432 the passive condition, the nonlocal models give similar results (Fig. 25).

433

434 Fig. 26 shows the strain localization pattern predicted by local model. The shear band  
 435 orientation in the backfill is directly measured from shear strain contours at  $u/h_w = 5\%$  (Fig.  
 436 27). The angle of the shear band under active earth condition is larger than that under passive  
 437 one. The angle of the shear band under active earth pressure decreases with increasing mesh  
 438 size, while that under passive earth pressure increases. For both cases, the angle range of the  
 439 local model ( $62^\circ - 66^\circ$ ) is larger than that of the nonlocal models ( $31^\circ - 36^\circ$ ). The nonlocal  
 440 functions reduce the range of measured angle which means they reduce mesh dependency,  
 441 especially for the G&S function, which is almost constant under active earth pressure.  
 442 Moreover, under passive earth pressure, the angle measured from the G&S function is slightly  
 443 larger than that of the GD and ON function.

444



445

446

447

448

**Fig. 24** The comparison of the retaining wall response on the sand under active failure condition: (a) Local model; (b) GD function; (c) G&S function; (d) ON function

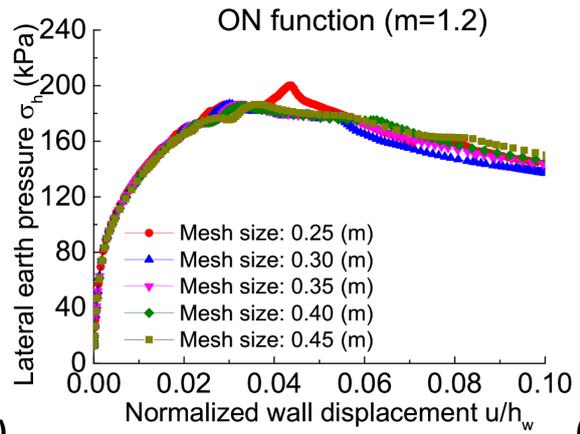
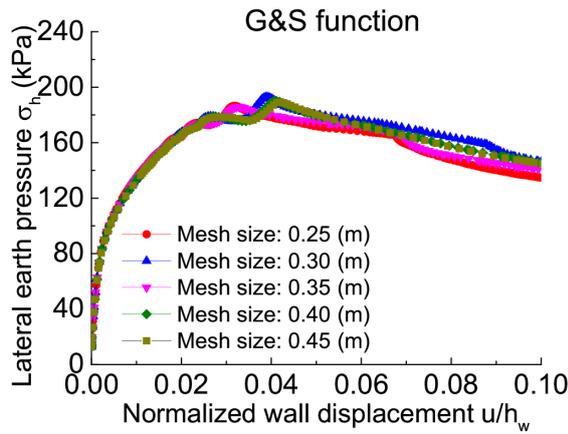
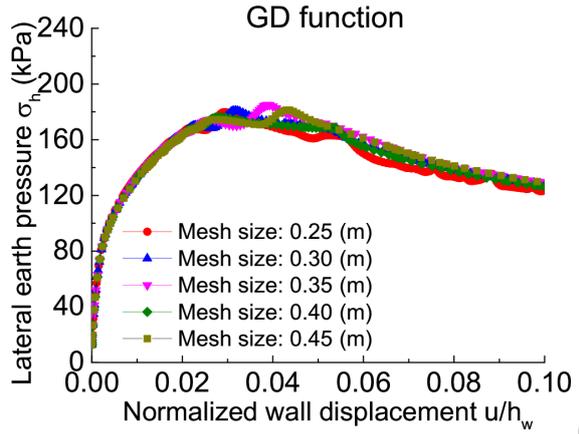
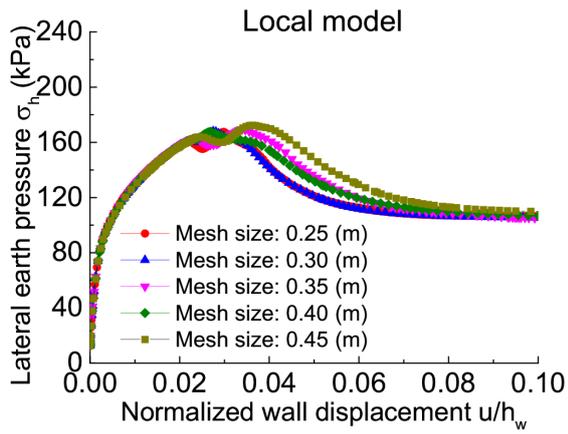
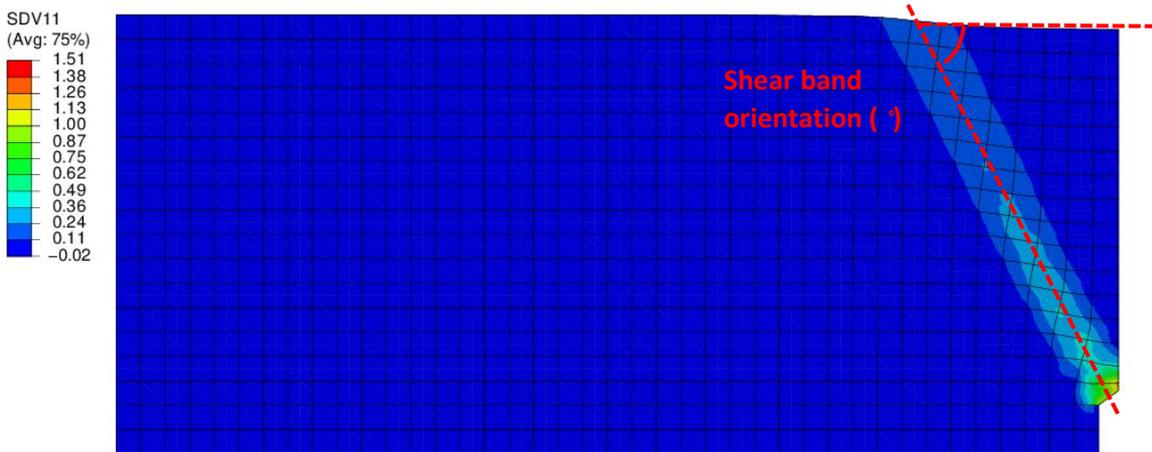
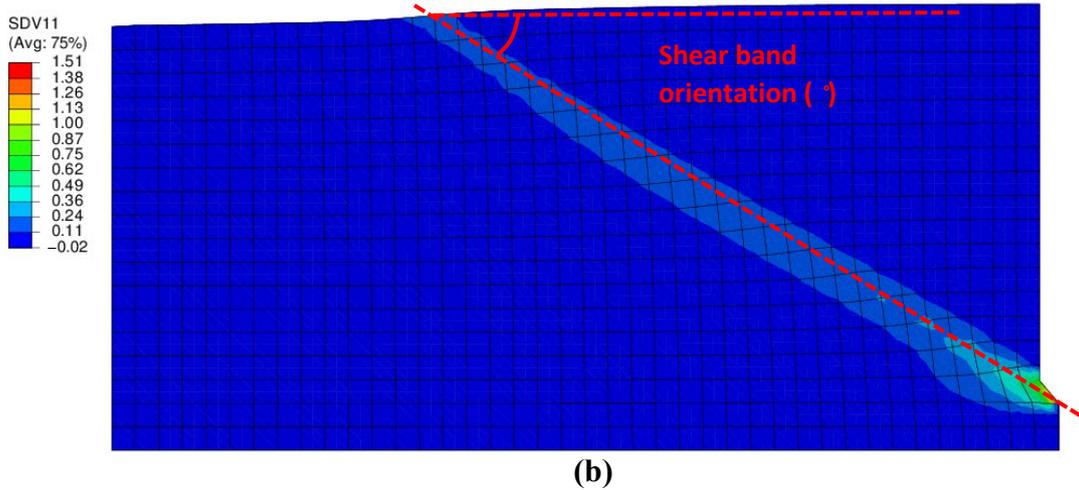


Fig. 25 Comparison of the retaining wall response on the sand under passive failure condition: (a) Local model; (b) GD function; (c) G&S function; (d) ON function



(a)



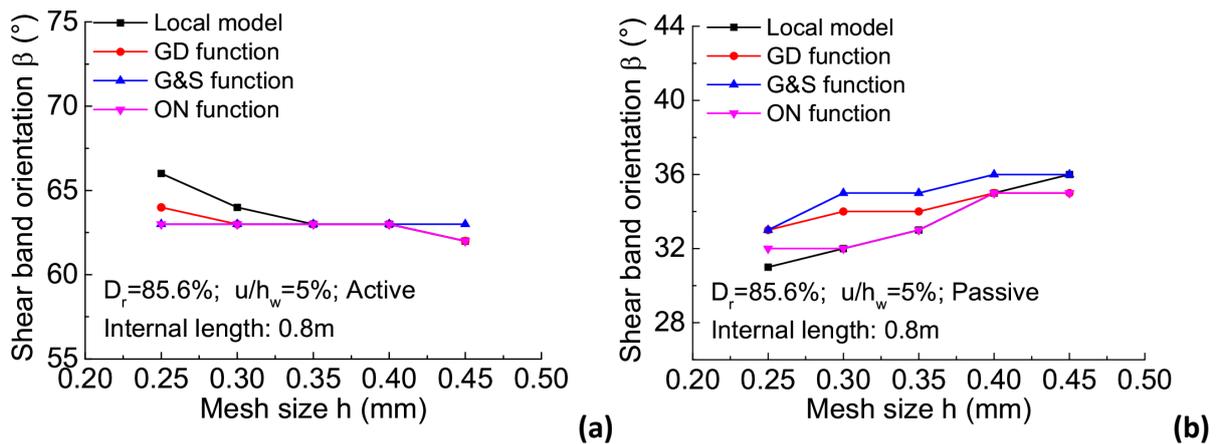
455

456 **Fig. 26 Shear band predicted by the G&S model after the retaining wall at  $u/h_w = 0.05$ :**

457

**(a) active condition and (b) passive condition**

458



459

460 **Fig. 27 Comparison of shear band orientation for retaining wall: (a) active failure condition**

461

**(b) passive failure condition**

462

## 463 7. Conclusion

464 The performance of three different weight functions for nonlocal regularisation have been  
 465 evaluated, including the GD, G&S and ON functions. An anisotropic sand model accounting for  
 466 evolution of anisotropy is used. The increment of void ratio is assumed nonlocal, which has a  
 467 significant influence on strain softening. Different BVPs have been simulated, including  
 468 drained and undrained plane strain compression, response of strip footings (level ground and  
 469 slope) and a retaining wall (passive and active conditions). The main conclusions are:

- 470 (a) All the nonlocal methods are effective in reducing the mesh dependency of the force-

471 displacement relationship in plane strain compression. The GD method gives less  
472 satisfactory results because the local value is contributing most to the nonlocal  
473 variable. The nonlocal regularization can reduce the mesh dependency of shear band  
474 thickness when the mesh size is smaller than the internal length. It is difficult to get  
475 mesh-independent shear band orientation in either drained or undrained condition.  
476 This could be due to that only the void ratio increment is assumed nonlocal. More  
477 mesh independent results could be obtained if more state variables that affect strain  
478 softening are assumed nonlocal.

479 (b) Nonlocal regularization can effectively reduce the mesh dependency of the force-  
480 displacement curves for strip footings. The ON method gives excessive overprediction  
481 of volume expansion for soil elements around the footings on a level ground, leading  
482 to an unrealistically steep reduction of the reaction force after peak.

483 (c) All three nonlocal functions give mesh-independent results for the active and passive  
484 earth pressures on the retaining wall. The shear band orientation predicted by the  
485 three functions shows small variation with the mesh size.

486 The G&S method is thus a better option for nonlocal regularisation of sand models. It does  
487 not require extra parameters and assumes that the local variable does not contribute to the  
488 nonlocal one. The GD function gives more mesh dependent results than the G&S function.  
489 The extra parameter  $m$  for the ON method can be determined using plane strain compression  
490 and used for the other BVPs. But the assumption that the local variable can make negative  
491 contribution to the nonlocal one may not be realistic. For instance, this assumption can cause  
492 very steep reduction of the reaction force on a strip footing on level sand ground.

493

#### 494 **Data Availability Statement**

495

496 All code, models, and data generated or used of this study are available from the  
497 corresponding author upon reasonable request.

498

#### 499 **Acknowledgements**

500

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503

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